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The quasi-linear Brezis-Nirenberg problem in low dimensions [☆]



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ABSTRACT

We discuss existence results for a quasi-linear elliptic equation of critical Sobolev growth [3,14] in the low-dimensional case, where the problem has a global character which is encoded in sign properties of the “regular” part for the corresponding Green’s function as in [9,11].

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1. Introduction

Let Ω be a bounded domain in \mathbb{R}^N , $N \geq 2$. Given $1 < p < N$ and $\lambda < \lambda_1$, let us discuss existence issues for the quasilinear problem

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$$\begin{cases} -\Delta_p u = \lambda u^{p-1} + u^{p^*-1} & \text{in } \Omega \\ u > 0 & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases} \tag{1.1}$$

where $\Delta_p(\cdot) = \operatorname{div}(|\nabla(\cdot)|^{p-2}\nabla(\cdot))$ is the p -Laplace operator, $p^* = \frac{Np}{N-p}$ is the so-called critical Sobolev exponent and λ_1 is the first eigenvalue of $-\Delta_p$ given by

$$\lambda_1 = \inf_{u \in W_0^{1,p}(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^p}{\int_{\Omega} |u|^p}.$$

Since $W_0^{1,p}(\Omega) \subset L^{p^*}(\Omega)$ is a continuous but non-compact embedding, standard variational methods fail to provide solutions of (1.1) by minimization of the Rayleigh quotient

$$Q_{\lambda}(u) = \frac{\int_{\Omega} |\nabla u|^p - \lambda \int_{\Omega} |u|^p}{\left(\int_{\Omega} |u|^{p^*}\right)^{\frac{p}{p^*}}}, \quad u \in W_0^{1,p}(\Omega) \setminus \{0\}.$$

Setting

$$S_{\lambda} = \inf \left\{ Q_{\lambda}(u) : u \in W_0^{1,p}(\Omega) \setminus \{0\} \right\},$$

it is known that S_0 coincides with the best Sobolev constant for the embedding $\mathcal{D}^{1,p}(\mathbb{R}^N) \subset L^{p^*}(\mathbb{R}^N)$ and then is never attained since independent of Ω . Moreover, by a Pohozaev identity $(1.1)_{\lambda=0}$ is not solvable on star-shaped domains, see [3,14]. The presence of the perturbation term λu^{p-1} in (1.1) can possibly restore compactness and produce minimizers for Q_{λ} , as shown for all $\lambda > 0$ first by Brezis and Nirenberg [3] in the semi-linear case when $N \geq 4$ and then by Guedda and Veron [14] when $N \geq p^2$.

Let us discuss now the low-dimensional case $p < N < p^2$. In the semi-linear situation $p = 2$ it corresponds to $N = 3$ and displays the following special features: according to [3], problem (1.1) is solvable on a ball precisely for $\lambda \in (\frac{\lambda_1}{4}, \lambda_1)$ and then, for the minimization problem on a general domain Ω , there holds

$$\lambda_* = \inf \{ \lambda \in (0, \lambda_1) : S_{\lambda} < S_0 \} \geq \frac{1}{4} \lambda_1(B) = \frac{\pi^2}{4} \left(\frac{3|\Omega|}{4\pi} \right)^{-\frac{2}{3}}$$

through a re-arrangement argument, where B is the ball having the same measure of Ω . In particular, for $\lambda \leq \frac{\lambda_1}{4}$ a general non-existence result on B follows from an integral identity of Pohozaev type, obtained by testing the equation against $\psi(|x|)u'$ for a suitable smooth function ψ with $\psi(0) = 0$. An integration by parts for the term

$$\int_0^1 r^{N-1} |u'|^{p-2} u' u \left[\frac{p-1}{p} \psi'' - \frac{N-1}{p} \frac{\psi'}{r} + \frac{N-1}{p} \frac{\psi}{r^2} \right]$$

is required to eliminate the dependence on the derivatives of u , which is possible in general just for $p = 2$. The property $\lambda^* > 0$ then requires a different proof for $p \neq 2$.

Since S_λ decreases in a continuous way from S_0 to 0 as λ ranges in $[0, \lambda_1)$, notice that $S_\lambda = S_0$ for $\lambda \in [0, \lambda_*]$, $S_\lambda < S_0$ for $\lambda \in (\lambda_*, \lambda_1)$ and S_λ is not attained for $\lambda \in [0, \lambda_*)$. A natural question concerns the case $\lambda = \lambda_*$ and the following general answer

$$S_{\lambda_*} \text{ is not achieved} \tag{1.2}$$

has been given by Druet [9], with an elegant proof which unfortunately seems not to work for $p \neq 2$. A complete characterization for the critical parameter λ_* then follows through a blow-up approach crucially based on (1.2).

We use here some of the results in [1] - precisely reported in Section 2 for reader's convenience - as a crucial ingredient to treat the quasilinear Brezis-Nirenberg problem (1.1) in the low-dimensional case $p < N < p^2$. Given $x_0 \in \Omega$ and $\lambda < \lambda_1$, introduce the Green function $G_\lambda(\cdot, x_0)$ as a positive solution to

$$\begin{cases} -\Delta_p G - \lambda G^{p-1} = \delta_{x_0} & \text{in } \Omega \\ G = 0 & \text{on } \partial\Omega. \end{cases} \tag{1.3}$$

Since uniqueness of $G_\lambda(\cdot, x_0)$ is just known for $p \geq 2$, hereafter we will just consider the case $p \geq 2$. If ω_N denotes the measure of the unit ball in \mathbb{R}^N , recall that the fundamental solution

$$\Gamma(x, x_0) = C_0 |x - x_0|^{-\frac{N-p}{p-1}}, \quad C_0 = \frac{p-1}{N-p} (N\omega_N)^{-\frac{1}{p-1}}, \tag{1.4}$$

solves $-\Delta_p \Gamma = \delta_{x_0}$ in \mathbb{R}^N . The function

$$H_\lambda(x, x_0) = G_\lambda(x, x_0) - \Gamma(x, x_0) \tag{1.5}$$

is usually referred to as the “regular” part of $G_\lambda(\cdot, x_0)$ but is just expected to be less singular than $\Gamma(x, x_0)$ at x_0 .

The complete characterization in [9] for λ_* (see also [11] for an alternative proof) still holds in the quasi-linear case, as stated by the following main result.

Theorem 1.1. *Let $2 \leq p < N < 2p$ and $0 < \lambda < \lambda_1$. The implications (i) \Rightarrow (ii) \Rightarrow (iii) do hold, where*

- (i) *there exists $x_0 \in \Omega$ such that $H_\lambda(x_0, x_0) > 0$*
- (ii) *$S_\lambda < S_0$*
- (iii) *S_λ is attained.*

Moreover, the implication (iii) \Rightarrow (i) does hold under the assumption (1.2) and in particular $\lambda_* > 0$.

Some comments are in order. Assumption $N < 2p$ is crucial here to guarantee that $H_\lambda(\cdot, x_0)$ is Hölder continuous at x_0 , see [1]. When $2p \leq N < p^2$ we conjecture $H_\lambda(x, x_0)$ to be mildly but still singular at x_0 , with a behavior like $\frac{m_\lambda(x_0)}{|x-x_0|^\alpha}$ for an appropriate $0 < \alpha < \frac{N-p}{p-1}$, and $m_\lambda(x_0)$ to play the same role as $H_\lambda(x_0, x_0)$ in Theorem 1.1. The quantity $m_\lambda(x_0)$ is usually referred to as the mass associated to $G_\lambda(\cdot, x_0)$ and appears in several contexts, see for example [12,13,18–20]. Notice that in the semilinear case $p = 2$ the range $2p \leq N < p^2$ is empty and such a situation doesn't show up in [9].

The implication $(iii) \Rightarrow (i)$ follows by a blow-up argument once (1.2) is assumed. To this aim, we first extend the pointwise blow-up theory in [10] to the quasi-linear context, a fundamental tool in the description of blow-up phenomena whose relevance goes beyond Theorem 1.1 and which completely settles some previous partial results [2,7,8] in this direction. Once sharp pointwise blow-up estimates are established, a major difficulty appears in the classical use of Pohozaev identities: written on small balls around the blow-up point as the radius tends to zero, they rule both the blow-up speed and the blow-up location since boundary terms in such identities can be controlled thanks to the property $\nabla H_\lambda(\cdot, x_0) \in L^\infty(\Omega)$. Clearly valid in the semi-linear situation, such gradient L^∞ -bound is completely missing in the quasi-linear context but surprisingly the correct answer can still be found by a different approach, based on a suitable approximation scheme for $G_\lambda(\cdot, x_0)$. At the same time, we provide a different proof of some facts in [9] in order to avoid some rough arguments concerning the limiting problems on halfspaces, when dealing with boundary blow-up.

Under the assumption (1.2), in the proof of Theorem 1.1 we will show that $H_{\lambda_*}(x_0, x_0) = 0$ for some $x_0 \in \Omega$, a stronger property than the validity of the implication $(iii) \Rightarrow (i)$ since $H_\lambda(x, x)$ is strictly increasing in λ for all $x \in \Omega$. Since S_0 is not attained, notice that (1.2) always holds if $\lambda_* = 0$ and then $\lambda_* > 0$ follows by the property $H_0(x_0, x_0) < 0$ for all $x_0 \in \Omega$. Moreover, since

$$\sup_{x \in \Omega} H_{\lambda_*}(x, x) = \max_{x \in \Omega} H_{\lambda_*}(x, x) = 0, \tag{1.6}$$

by monotonicity of H_λ in λ and under the assumption (1.2) the critical parameter λ_* is the first unique value of $\lambda > 0$ attaining (1.6) and can be re-written as

$$\lambda_* = \sup \{ \lambda \in (0, \lambda_1) : H_\lambda(x, x) < 0 \text{ for all } x \in \Omega \} .$$

In Section 2 we recall some facts from [1] that will be used throughout the paper and prove some useful convergence properties. The implication $(i) \Rightarrow (ii)$ is established in Section 3 by the expansion of $Q_\lambda(PU_{\epsilon, x_0})$ along the “bubble” PU_{ϵ, x_0} concentrating at x_0 as $\epsilon \rightarrow 0$ and integral identities of Pohozaev type for $G_\lambda(\cdot, x_0)$, crucial for a fine asymptotic analysis, are also derived. Section 4 is devoted to develop the blow-up argument along with sharp pointwise estimates to establish the final part in Theorem 1.1.

2. Some preliminary facts

For reader's convenience, let us collect here some of the results in [1]. To give the statement of Theorem 1.1 a full meaning, we need a general theory for problem (1.3), as stated in the following result.

Theorem 2.1. [1] *Let $1 < p \leq N$ and $\lambda < \lambda_1$. Assume $p \geq 2$ and $N < 2p$ if $\lambda \neq 0$. Then problem (1.3) has a positive solution $G_\lambda(\cdot, x_0)$ so that $H_\lambda(x, x_0)$ in (1.5) satisfies*

$$\nabla H_\lambda(\cdot, x_0) \in L^{\bar{q}}(\Omega), \quad \bar{q} = \frac{N(p-1)}{N-1}, \tag{2.1}$$

which is unique when either $\lambda = 0$ or $\lambda \neq 0$ and (2.1) holds. Moreover

- given $M > 0$, $q_0 > \frac{N}{p}$ and $p_0 \geq 1$ there exists $C > 0$ so that

$$\|H + c\|_{\infty, B_r(x_0)} \leq C(r^{-\frac{N}{p_0}} \|H + c\|_{p_0, B_{2r}(x_0)} + r^{\frac{pq_0 - N}{q_0(p-1)}} \|f\|_{\frac{1}{p-1}, q_0, B_{2r}(x_0)}) \tag{2.2}$$

for all $\epsilon, r, c \in \mathbb{R}$, $f \in L^{q_0}(\Omega)$ and solution $G = \Gamma + H$, with $H \in L^\infty(\Omega)$ and $\nabla H \in L^{\bar{q}}(\Omega)$, to

$$-\Delta_p G + \Delta_p \Gamma = f \quad \text{in } \Omega \setminus \{x_0\} \tag{2.3}$$

so that $\epsilon^{p-1} \leq r \leq \frac{1}{4} \text{dist}(x_0, \partial\Omega)$, $\frac{|x-x_0|^{\frac{1}{p-1}}}{M(\epsilon^p + |x-x_0|^{\frac{p}{p-1}})^{\frac{N}{p}}} \leq |\nabla \Gamma| \leq M|\nabla \Gamma|(x, x_0)$, $|c| +$

$\|H\|_\infty + \|f\|_{\frac{1}{q_0}, \frac{1}{p-1}} \leq M$, where $\Gamma(\cdot, x_0)$ is given by (1.4);

- $\lambda G_\lambda^{p-1} \in L^{q_0}(\Omega)$ for $q_0 > \frac{N}{p}$ and $H_\lambda(\cdot, x_0)$ is a continuous function in $\bar{\Omega}$ satisfying

$$|H_\lambda(x, x_0) - H_\lambda(x_0, x_0)| \leq C|x - x_0|^\alpha \quad \forall x \in \Omega \tag{2.4}$$

for some $C > 0$, $\alpha \in (0, 1)$ with $H_\lambda(x_0, x_0)$ strictly increasing in λ .

Notice that the first part in Theorem 2.1 has been established in [15]. Let us stress that the condition $f \in L^{q_0}(\Omega)$ for some $q_0 > \frac{N}{p}$, which is valid for $f = \lambda G_\lambda^{p-1}$ when $N < 2p$ if $\lambda \neq 0$, is a natural condition on the R.H.S. of the difference equation (2.3) to prove L^∞ -bounds on H as it arises for instance in the Moser iterative argument adopted in [22]. In this respect, observe that also in the semilinear case $H_\lambda(\cdot, x_0)$ is no longer regular at x_0 when $4 = 2p \leq N$.

The following a-priori estimates are the basis of Theorem 2.1 and will be crucially used here to establish some accurate pointwise blow-up estimates.

Proposition 2.2. [1] *Let $2 \leq p \leq N$. Assume that $a_n \in L^\infty(\Omega)$, $f_n \in L^1(\Omega)$ and g_n, \hat{g}_n satisfy*

$g_n, \hat{g}_n \in L^\infty(\Omega) \cap W^{1,p}(\Omega)$ p -harmonic in Ω , g_n, \hat{g}_n non-constant unless 0

and

$$\lim_{n \rightarrow +\infty} \|a_n - a\|_\infty = 0 \text{ with } \sup_\Omega a < \lambda_1, \quad \sup_{n \in \mathbb{N}} [\|f_n\|_1 + \|g_n\|_\infty + \|\hat{g}_n\|_\infty] < +\infty.$$

If $u_n \in W_{g_n}^{1,p}(\Omega)$ solves $-\Delta_p u_n - a_n |u_n|^{p-2} u_n = f_n$ in Ω , then $\sup_{n \in \mathbb{N}} \|u_n\|_{p-1} < +\infty$ and, if $g_n = g$, the sequence u_n is pre-compact in $W^{1,q}(\Omega)$ for all $1 \leq q < \bar{q}$. Moreover, if $N < 2p$, $a_n = \lambda_n \in \mathbb{R}$ and $\hat{u}_n \in W_{\hat{g}_n}^{1,p}(\Omega)$ solves $-\Delta_p \hat{u}_n = f_n$ in Ω , then $\sup_{n \in \mathbb{N}} \|u_n - \hat{u}_n\|_\infty < \infty$.

We will also make use of the following general form of comparison principle.

Proposition 2.3. [1] Let $2 \leq p \leq N$ and $a, f_1, f_2 \in L^\infty(\Omega)$. Let $u_i \in C^1(\bar{\Omega})$, $i = 1, 2$, be solutions to

$$-\Delta_p u_i - a u_i^{p-1} = f_i \quad \text{in } \Omega$$

so that

$$u_i > 0 \text{ in } \Omega, \quad \frac{u_1}{u_2} \leq C \text{ near } \partial\Omega$$

for some $C > 0$. If $f_1 \leq f_2$ with $f_2 \geq 0$ in Ω and $u_1 \leq u_2$ on $\partial\Omega$, then $u_1 \leq u_2$ in Ω .

Let us introduce now a special approximation scheme for $G_\lambda(\cdot, x_0)$, which is particularly suited for the problem we are interested in. Given $C_1 = N^{\frac{N-p}{p^2}} \left(\frac{N-p}{p-1}\right)^{\frac{(p-1)(N-p)}{p^2}}$, the so-called standard bubbles

$$U_{\epsilon, x_0}(x) = C_1 \left(\frac{\epsilon}{\epsilon^p + |x - x_0|^{\frac{p}{p-1}}} \right)^{\frac{N-p}{p}} \quad \epsilon > 0, \quad x_0 \in \mathbb{R}^N, \tag{2.5}$$

are the extremals of the Sobolev inequality

$$S_0 \left(\int_{\mathbb{R}^N} |u|^{p^*} \right)^{\frac{p}{p^*}} \leq \int_{\mathbb{R}^N} |\nabla u|^p, \quad u \in \mathcal{D}^{1,p}(\mathbb{R}^N),$$

and the unique entire solutions in $\mathcal{D}^{1,p}(\mathbb{R}^N)$ of

$$-\Delta_p U = U^{p^*-1} \quad \text{in } \mathbb{R}^N, \tag{2.6}$$

see [5,21,25]. For $\lambda < \lambda_1$ consider its projection PU_{ϵ, x_0} in Ω , as the solution of

$$\begin{cases} -\Delta_p PU_{\epsilon,x_0} = \lambda PU_{\epsilon,x_0}^{p-1} + U_{\epsilon,x_0}^{p^*-1} & \text{in } \Omega \\ PU_{\epsilon,x_0} > 0 & \text{in } \Omega \\ PU_{\epsilon,x_0} = 0 & \text{on } \partial\Omega. \end{cases} \tag{2.7}$$

Letting $G_{\epsilon,x_0} = \frac{C_0}{C_1} \epsilon^{-\frac{N-p}{p}} PU_{\epsilon,x_0}$ with C_0 given by (1.4), decompose it as $G_{\epsilon,x_0} = \Gamma_{\epsilon,x_0} + H_{\epsilon,x_0}$, where

$$\Gamma_{\epsilon,x_0} = \frac{C_0}{C_1} \epsilon^{-\frac{N-p}{p}} U_{\epsilon,x_0} = \frac{C_0}{(\epsilon^p + |x - x_0|^{\frac{p}{p-1}})^{\frac{N-p}{p}}} \rightarrow \Gamma(x, x_0) \tag{2.8}$$

in $C^1_{\text{loc}}(\bar{\Omega} \setminus \{x_0\})$ as $\epsilon \rightarrow 0$. Since

$$f_{\epsilon,x_0} := -\Delta_p \Gamma_{\epsilon,x_0} = \left(\frac{C_0}{C_1} \epsilon^{-\frac{N-p}{p}}\right)^{p-1} U_{\epsilon,x_0}^{p^*-1} = \frac{C_0^{p-1} C_1^{\frac{p^2}{N-p}} \epsilon^p}{(\epsilon^p + |x - x_0|^{\frac{p}{p-1}})^{N-\frac{N-p}{p}}} \rightarrow 0 \tag{2.9}$$

in $C^1_{\text{loc}}(\bar{\Omega} \setminus \{x_0\})$ and

$$\int_{\Omega} f_{\epsilon,x_0} = - \int_{\partial\Omega} |\nabla \Gamma_{\epsilon,x_0}|^{p-2} \partial_{\nu} \Gamma_{\epsilon,x_0} \rightarrow - \int_{\partial\Omega} |\nabla \Gamma|^{p-2}(x, x_0) \partial_{\nu} \Gamma(x, x_0) d\sigma(x) = 1$$

as $\epsilon \rightarrow 0$ in view of (2.6) and (2.8), notice that $f_{\epsilon,x_0} \rightharpoonup \delta_{x_0}$ weakly in the sense of measures in Ω as $\epsilon \rightarrow 0$ and G_{ϵ,x_0} solves

$$\begin{cases} -\Delta_p G_{\epsilon,x_0} = \lambda G_{\epsilon,x_0}^{p-1} + f_{\epsilon,x_0} & \text{in } \Omega \\ G_{\epsilon,x_0} > 0 & \text{in } \Omega \\ G_{\epsilon,x_0} = 0 & \text{on } \partial\Omega. \end{cases} \tag{2.10}$$

Thanks to Theorem 2.1 and Proposition 2.2 we can now establish the following convergence result.

Proposition 2.4. *Let $2 \leq p \leq N$ and assume $N < 2p$ if $\lambda \neq 0$. Then there holds*

$$H_{\epsilon,x_0} \rightarrow H_{\lambda}(\cdot, x_0) \quad \text{in } C(\bar{\Omega}) \tag{2.11}$$

as $\epsilon \rightarrow 0$.

Proof. By Proposition 2.2 we can find a subsequence $\epsilon_n \rightarrow 0$ so that $G_{\epsilon_n,x_0} \rightarrow G$ in $W^{1,q}_0(\Omega)$ as $n \rightarrow +\infty$ for all $1 \leq q < \bar{q}$, where $G = \Gamma(x, x_0) + H$ is a solution of (1.3) for some H in view of (2.8) and (2.10). In particular, if $\lambda \neq 0$ by the Sobolev embedding theorem there holds

$$G_{\epsilon_n,x_0} \rightarrow G \quad \text{in } L^p(\Omega) \text{ as } n \rightarrow +\infty \tag{2.12}$$

thanks to $\bar{q}^* > p$ in view of $N < 2p \leq p^2$. Moreover, let us rewrite (2.10) in the equivalent form:

$$\begin{cases} -\Delta_p(\Gamma_{\epsilon,x_0} + H_{\epsilon,x_0}) + \Delta_p \Gamma_{\epsilon,x_0} = \lambda G_{\epsilon,x_0}^{p-1} & \text{in } \Omega \\ H_{\epsilon,x_0} = -\Gamma_{\epsilon,x_0} & \text{on } \partial\Omega. \end{cases} \tag{2.13}$$

Let us denote the solution of (2.10) $_{\lambda=0}$ by G_{ϵ,x_0}^0 and set $H_{\epsilon,x_0}^0 = G_{\epsilon,x_0}^0 - \Gamma_{\epsilon,x_0}$. By the uniqueness part in Theorem 2.1 with $\lambda = 0$ we have that

$$G_{\epsilon,x_0}^0 \rightarrow G_0(\cdot, x_0) \text{ in } W_0^{1,q}(\Omega)$$

as $\epsilon \rightarrow 0$, for all $1 \leq q < \bar{q}$. Moreover, since $|H_{\epsilon,x_0}^0| \leq M$ on $\partial\Omega$, by integrating (2.13) against $(H_{\epsilon,x_0}^0 \mp M)_{\pm}$ we deduce that

$$|H_{\epsilon,x_0}^0| \leq M \quad \text{in } \Omega \tag{2.14}$$

in an uniform way and then G_{ϵ,x_0}^0 is locally uniformly bounded in $\bar{\Omega} \setminus \{x_0\}$. By elliptic estimates [6,16,22,23] and (2.10) $_{\lambda=0}$ we deduce that

$$G_{\epsilon,x_0}^0 \text{ uniformly bounded in } C_{\text{loc}}^{1,\alpha}(\bar{\Omega} \setminus \{x_0\}) \tag{2.15}$$

for some $\alpha \in (0, 1)$. Integrating (2.13) $_{\lambda=0}$ against $\eta^p H_{\epsilon,x_0}^0$, $0 \leq \eta \in C_0^\infty(\Omega)$, we get that

$$\int_{\Omega} \eta^p |\nabla H_{\epsilon,x_0}^0|^p \leq p \int_{\Omega} \eta^{p-1} |\nabla \eta| (|\nabla \Gamma_{\epsilon,x_0}|^{p-2} + |\nabla H_{\epsilon,x_0}^0|^{p-2}) |H_{\epsilon,x_0}^0| |\nabla H_{\epsilon,x_0}^0|$$

and then (2.14) and Young’s inequality imply that

$$\nabla H_{\epsilon,x_0}^0 \text{ uniformly bounded in } L^p(\Omega) \tag{2.16}$$

in view of (2.15).

Let us consider now the case $\lambda \neq 0$. Since

$$-\Delta_p(\Gamma_{\epsilon,x_0} + H_{\epsilon,x_0}) + \Delta_p(\Gamma_{\epsilon,x_0} + H_{\epsilon,x_0}^0) = \lambda G_{\epsilon,x_0}^{p-1} \quad \text{in } \Omega$$

with $H_{\epsilon,x_0} - H_{\epsilon,x_0}^0 = 0$ on $\partial\Omega$, an integration against $H_{\epsilon,x_0} - H_{\epsilon,x_0}^0$ gives that

$$\int_{\Omega} |\nabla (H_{\epsilon,x_0} - H_{\epsilon,x_0}^0)|^p \leq |\lambda| \int_{\Omega} G_{\epsilon,x_0}^{p-1} |H_{\epsilon,x_0} - H_{\epsilon,x_0}^0| \leq |\lambda| \|G_{\epsilon,x_0}\|_p^{p-1} \|H_{\epsilon,x_0} - H_{\epsilon,x_0}^0\|_p$$

thanks to the Hölder’s inequality and the coercivity properties of the p -Laplace operator, and then

$$\nabla (H_{\epsilon_n, x_0}^0 - H_{\epsilon_n, x_0}^0) \text{ uniformly bounded in } L^p(\Omega) \tag{2.17}$$

in view of (2.12) and Poincaré inequality. A combination of (2.16) and (2.17) lead to a uniform L^p -bound on $\nabla H_{\epsilon_n, x_0}^0$, showing by Fatou’s lemma that $\nabla H \in L^p(\Omega)$. By Theorem 2.1 we have that $G = G_\lambda(\cdot, x_0)$ and then

$$G_{\epsilon, x_0} \rightarrow G_\lambda(\cdot, x_0) \text{ in } W_0^{1,q}(\Omega) \tag{2.18}$$

as $\epsilon \rightarrow 0$, for all $1 \leq q < \bar{q}$.

To extend (2.14) to the case $\lambda \neq 0$, observe that (2.10) and $-\Delta_p \Gamma_{\epsilon, x_0} = f_{\epsilon, x_0}$ in Ω imply $\|H_{\epsilon, x_0}\|_\infty \leq C$ for all $\epsilon > 0$ thanks to Proposition 2.2 in view of $N < 2p$ when $\lambda \neq 0$. Since $f = \lambda G_{\epsilon, x_0}^{p-1}$ is uniformly bounded in $L^{q_0}(\Omega)$ for some $q_0 > \frac{N}{p}$ in view of $\frac{\bar{q}^*}{p-1} > \frac{N}{p}$ when $N < 2p$ and

$$|\nabla \Gamma_{\epsilon, x_0}| = \frac{C_0(N-p)}{p-1} \frac{|x-x_0|^{\frac{1}{p-1}}}{(\epsilon^p + |x-x_0|^{\frac{p}{p-1}})^{\frac{N}{p}}} \leq M |\nabla \Gamma|(x, x_0),$$

we can apply (2.2) in Theorem 2.1 to H_{ϵ, x_0} as a solution to (2.13) by getting

$$|H_{\epsilon, x_0}(x) - H_\lambda(x_0, x_0)| \leq C \left(r^{-\frac{N}{p-1}} \|H_{\epsilon, x_0} - H_\lambda(x_0, x_0)\|_{p-1, B_{2r}(x_0)} + r^{\frac{pq_0-N}{q_0(p-1)}} \right) \tag{2.19}$$

for all $x \in B_r(x_0)$ and $\epsilon^{p-1} \leq r \leq \frac{1}{4} \text{dist}(x_0, \partial\Omega)$.

By contradiction assume that (2.11) does not hold. Then there exist sequences $\epsilon_n \rightarrow 0$ and $x_n \in \Omega$ so that $|H_{\epsilon_n, x_0}(x_n) - H_\lambda(x_n, x_0)| \geq 2\delta > 0$. Since by elliptic estimates [6,16,22,23] there holds

$$G_{\epsilon, x_0} \rightarrow G_\lambda(\cdot, x_0) \text{ in } C_{\text{loc}}^1(\bar{\Omega} \setminus \{x_0\}) \tag{2.20}$$

as $\epsilon \rightarrow 0$ in view of (2.10) and (2.18), we have that $\bar{x} = x_0$ and then

$$|H_{\epsilon_n, x_0}(x_n) - H_\lambda(x_0, x_0)| \geq \delta \tag{2.21}$$

thanks to $H_\lambda(\cdot, x_0) \in C(\bar{\Omega})$. Since by the Sobolev embedding theorem $H_{\epsilon, x_0} \rightarrow H_\lambda(\cdot, x_0)$ in $L^{p-1}(\Omega)$ as $\epsilon \rightarrow 0$ in view of (2.18) and $\bar{q}^* > p-1$, we can insert (2.21) into (2.19) and get as $n \rightarrow +\infty$

$$\delta \leq C \left(r^{-\frac{N}{p-1}} \|H_\lambda(\cdot, x_0) - H_\lambda(x_0, x_0)\|_{p-1, B_{2r}(x_0)} + r^{\frac{pq_0-N}{q_0(p-1)}} \right) \tag{2.22}$$

for all $0 < r \leq \frac{1}{4} \text{dist}(x_0, \partial\Omega)$. Since

$$r^{-\frac{N}{p-1}} \|H_\lambda(\cdot, x_0) - H_\lambda(x_0, x_0)\|_{p-1, B_{2r}(x_0)} \leq Cr^\alpha \rightarrow 0$$

as $r \rightarrow 0$ thanks to (2.4), estimate (2.22) leads to a contradiction and the proof is complete. \square

As a by-product we have the following useful result.

Corollary 2.5. *Let $2 \leq p \leq N$ and assume $N < 2p$ if $\lambda \neq 0$. Then the expansion*

$$PU_{\epsilon, x_0} = U_{\epsilon, x_0} + \frac{C_1}{C_0} \epsilon^{\frac{N-p}{p}} H_\lambda(\cdot, x_0) + o\left(\epsilon^{\frac{N-p}{p}}\right) \tag{2.23}$$

does hold uniformly in Ω as $\epsilon \rightarrow 0$.

3. Energy expansions and Pohozaev identities

We are concerned with the discussion of implication (i) \Rightarrow (ii) in Theorem 1.1, whereas the proof of (ii) \Rightarrow (iii) in Theorem 1.1 is rather classical and can be found in [14].

Let $0 < \lambda < \lambda_1$ and $x_0 \in \Omega$ so that $H_\lambda(x_0, x_0) > 0$. In order to show $S_\lambda < S_0$ let us expand $Q_\lambda(PU_{\epsilon, x_0})$ for $\epsilon > 0$ small. Since PU_{ϵ, x_0} solves (2.7), we have that

$$\begin{aligned} \int_{\Omega} |\nabla PU_{\epsilon, x_0}|^p - \lambda \int_{\Omega} (PU_{\epsilon, x_0})^p &= \int_{\Omega} U_{\epsilon, x_0}^{p^*-1} PU_{\epsilon, x_0} = \int_{\Omega} U_{\epsilon, x_0}^{p^*} \\ &+ \frac{C_1}{C_0} \epsilon^{\frac{N-p}{p}} \int_{\Omega} U_{\epsilon, x_0}^{p^*-1} [H_\lambda(x, x_0) + o(1)] \end{aligned} \tag{3.1}$$

as $\epsilon \rightarrow 0$ in view of (2.23). Given $\Omega_\epsilon = \frac{\Omega - x_0}{\epsilon^{p-1}}$ observe that

$$\int_{\Omega} U_{\epsilon, x_0}^{p^*} = \int_{\Omega_\epsilon} U_1^{p^*} = \int_{\mathbb{R}^N} U_1^{p^*} + O(\epsilon^N) \tag{3.2}$$

and

$$\begin{aligned} \int_{\Omega} U_{\epsilon, x_0}^{p^*-1} [H_\lambda(x, x_0) + o(1)] &= \int_{\Omega} U_{\epsilon, x_0}^{p^*-1} [H_\lambda(x_0, x_0) + O(|x - x_0|^\alpha) + o(1)] \\ &= \epsilon^{\frac{(N-p)(p-1)}{p}} \int_{\Omega_\epsilon} U_1^{p^*-1} [H_\lambda(x_0, x_0) + O(\epsilon^{\alpha(p-1)} |y|^\alpha) + o(1)] \\ &= \epsilon^{\frac{(N-p)(p-1)}{p}} H_\lambda(x_0, x_0) \int_{\mathbb{R}^N} U_1^{p^*-1} + o\left(\epsilon^{\frac{(N-p)(p-1)}{p}}\right) \end{aligned} \tag{3.3}$$

in view of (2.4) and $\int_{\mathbb{R}^n} U_1^{p^*-1} |y|^\alpha < +\infty$. Inserting (3.2)-(3.3) into (3.1) we deduce

$$\int_{\Omega} |\nabla PU_{\epsilon, x_0}|^p - \lambda \int_{\Omega} (PU_{\epsilon, x_0})^p = \int_{\mathbb{R}^N} U_1^{p^*} + \epsilon^{N-p} \frac{C_1}{C_0} H_\lambda(x_0, x_0) \int_{\mathbb{R}^N} U_1^{p^*-1} + o(\epsilon^{N-p}). \tag{3.4}$$

By the Taylor expansion

$$(PU_{\epsilon, x_0})^{p^*} = U_{\epsilon, x_0}^{p^*} + \epsilon^{\frac{N-p}{p}} \frac{C_1}{C_0} p^* U_{\epsilon, x_0}^{p^*-1} [H_\lambda(x, x_0) + o(1)] + O(\epsilon^{2\frac{N-p}{p}} U_{\epsilon, x_0}^{p^*-2} + \epsilon^N)$$

in view of (2.23) and $\|H_\lambda(\cdot, x_0)\|_\infty < +\infty$, we obtain

$$\int_{\Omega} (PU_{\epsilon, x_0})^{p^*} = \int_{\mathbb{R}^N} U_1^{p^*} + \epsilon^{N-p} \frac{C_1}{C_0} p^* H_\lambda(x_0, x_0) \int_{\mathbb{R}^N} U_1^{p^*-1} + o(\epsilon^{N-p}) \tag{3.5}$$

thanks to (3.2)-(3.3) and

$$\int_{\Omega} U_{\epsilon, x_0}^{p^*-2} = \epsilon^{2\frac{(N-p)(p-1)}{p}} \int_{\Omega_\epsilon} U_1^{p^*-2} = O(\epsilon^{2\frac{(N-p)(p-1)}{p}})$$

for $N < 2p$. Expansions (3.4)-(3.5) now yield

$$Q_\lambda(PU_{\epsilon, x_0}) = S_0 - (p-1) S_0^{\frac{p-N}{p}} \left(\int_{\mathbb{R}^N} U_1^{p^*-1} \right) \frac{C_1}{C_0} \epsilon^{N-p} H_\lambda(x_0, x_0) + o(\epsilon^{N-p})$$

in view of (2.6) and

$$S_0 = \frac{\int_{\mathbb{R}^N} |\nabla U_1|^p}{\left(\int_{\mathbb{R}^N} U_1^{p^*}\right)^{\frac{p}{p^*}}} = \left(\int_{\mathbb{R}^N} U_1^{p^*}\right)^{\frac{p}{N}}.$$

Then, for $\epsilon > 0$ small we obtain that $S_\lambda < S_0$ thanks to $H_\lambda(x_0, x_0) > 0$.

As already discussed in the Introduction, a fundamental tool is represented by the Pohozaev identity. Derived [4] for autonomous PDE's involving the p -Laplace operator, it extends to the non-autonomous case and writes, in the situation of our interest, as follows: if $u \in C^{1,\alpha}(\bar{D})$ solves $-\Delta_p u = \lambda u^{p-1} + cu^{p^*-1} + f$ in D for $f \in C^1(\bar{D})$ and $c \in \{0, 1\}$, given $x_0 \in \mathbb{R}^N$ there holds

$$\begin{aligned} & \int_D [NH - f\langle x - x_0, \nabla u \rangle - \frac{N-p}{p} |\nabla u|^p] \\ &= \int_{\partial D} \langle x - x_0, -\frac{|\nabla u|^p}{p} \nu + |\nabla u|^{p-2} \partial_\nu u \nabla u + H\nu \rangle \end{aligned} \tag{3.6}$$

with $H(u) = \frac{\lambda}{p} u^p + \frac{c}{p^*} u^{p^*}$ and

$$\int_D |\nabla u|^p = \int_D [\lambda u^p + cu^{p^*} + fu] + \int_{\partial D} u |\nabla u|^{p-2} \partial_\nu u. \tag{3.7}$$

An integral identity of Pohozaev type for $G_\lambda(\cdot, x_0)$ like (3.8) below is of fundamental importance since $H_\lambda(x_0, x_0)$ appears as a sort of residue. In the semi-linear case

such identity (3.8) holds in the limit of (3.6)-(3.7) on $B_\delta(x_0)$ as $\delta \rightarrow 0$ thanks to $\nabla H_\lambda(\cdot, x_0) \in L^\infty(\Omega)$, a property far from being obvious in the quasi-linear context where just integral bounds on $\nabla H_\lambda(\cdot, x_0)$ like (2.1) are available. Instead, we can use the special approximating sequence G_{ϵ, x_0} to derive the following result.

Proposition 3.1. *Let $2 \leq p < N$ and assume $N < 2p$ if $\lambda \neq 0$. Given $x_0 \in \Omega$, $0 < \delta < \text{dist}(x_0, \partial\Omega)$ and $\lambda < \lambda_1$, there holds*

$$\begin{aligned}
 & C_0 H_\lambda(x_0, x_0) \\
 &= \lambda \int_{B_\delta(x_0)} G_\lambda^p(x, x_0) dx + \int_{\partial B_\delta(x_0)} \left(\frac{\delta}{p} |\nabla G_\lambda(x, x_0)|^p - \delta |\nabla G_\lambda(x, x_0)|^{p-2} (\partial_\nu G_\lambda(x, x_0))^2 \right. \\
 &\quad \left. - \frac{\lambda \delta}{p} G_\lambda^p(x, x_0) - \frac{N-p}{p} G_\lambda(x, x_0) |\nabla G_\lambda(x, x_0)|^{p-2} \partial_\nu G_\lambda(x, x_0) \right) d\sigma(x) \tag{3.8}
 \end{aligned}$$

for some $C_0 > 0$.

Proof. Since by elliptic regularity theory [6,16,22,23] $G_{\epsilon, x_0} \in C^{1,\alpha}(\bar{\Omega})$ for some $\alpha \in (0, 1)$ in view of (2.10), we can apply the Pohozaev identity (3.6) to G_{ϵ, x_0} with $c = 0$ and $f = f_{\epsilon, x_0}$ on $D = B_\delta(x_0) \subset \Omega$ to get

$$\begin{aligned}
 & \int_{\partial B_\delta(x_0)} \left(-\frac{\delta}{p} |\nabla G_{\epsilon, x_0}|^p + \delta |\nabla G_{\epsilon, x_0}|^{p-2} (\partial_\nu G_{\epsilon, x_0})^2 + \frac{\lambda \delta}{p} G_{\epsilon, x_0}^p \right. \\
 &\quad \left. + \frac{N-p}{p} G_{\epsilon, x_0} |\nabla G_{\epsilon, x_0}|^{p-2} \partial_\nu G_{\epsilon, x_0} \right) \\
 &= \int_{B_\delta(x_0)} \left(\lambda G_{\epsilon, x_0}^p - \frac{N-p}{p} f_{\epsilon, x_0} G_{\epsilon, x_0} - f_{\epsilon, x_0} \langle x - x_0, \nabla G_{\epsilon, x_0} \rangle \right) \tag{3.9}
 \end{aligned}$$

in view of (3.7). The approximating sequence G_{ϵ, x_0} has the key property that $\nabla G_{\epsilon, x_0}$ and f_{ϵ, x_0} are at main order multiples of $\nabla U_{\epsilon, x_0}$ and $U_{\epsilon, x_0}^{p^*-1}$, respectively, in such a way that $f_{\epsilon, x_0} \nabla G_{\epsilon, x_0}$ allows for a further integration by parts of the R.H.S. in (3.9). The function H_{ϵ, x_0} appears in the remaining lower-order terms and explains why in the limit $\epsilon \rightarrow 0$ an additional term containing $H_\lambda(x_0, x_0)$ will appear in (3.8). The identity

$$\begin{aligned}
 & \int_{B_\delta(x_0)} U_{\epsilon, x_0}^{p^*-2} G_{\epsilon, x_0} \langle x - x_0, \nabla U_{\epsilon, x_0} \rangle \\
 &= \int_{B_\delta(x_0)} U_{\epsilon, x_0}^{p^*-1} \langle x - x_0, \nabla G_{\epsilon, x_0} - \nabla H_{\epsilon, x_0} + H_{\epsilon, x_0} \frac{\nabla U_{\epsilon, x_0}}{U_{\epsilon, x_0}} \rangle \\
 &= \int_{B_\delta(x_0)} U_{\epsilon, x_0}^{p^*-1} \langle x - x_0, \nabla G_{\epsilon, x_0} + p^* H_{\epsilon, x_0} \frac{\nabla U_{\epsilon, x_0}}{U_{\epsilon, x_0}} \rangle
 \end{aligned}$$

$$-\delta \int_{\partial B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-1} H_{\epsilon,x_0} + N \int_{B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-1} H_{\epsilon,x_0}$$

does hold thanks to $G_{\epsilon,x_0} = \Gamma_{\epsilon,x_0} + H_{\epsilon,x_0}$ and $\Gamma_{\epsilon,x_0} \nabla U_{\epsilon,x_0} = U_{\epsilon,x_0} (\nabla G_{\epsilon,x_0} - \nabla H_{\epsilon,x_0})$, which inserted into

$$\begin{aligned} & \int_{B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-1} \langle x - x_0, \nabla G_{\epsilon,x_0} \rangle \\ &= \delta \int_{\partial B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-1} G_{\epsilon,x_0} - (p^* - 1) \int_{B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-2} G_{\epsilon,x_0} \langle x - x_0, \nabla U_{\epsilon,x_0} \rangle \\ & \quad - N \int_{B_\delta(x_0)} U_{\epsilon,x_0}^{p^*-1} G_{\epsilon,x_0} \end{aligned}$$

leads to

$$\begin{aligned} \int_{B_\delta(x_0)} f_{\epsilon,x_0} \langle x - x_0, \nabla G_{\epsilon,x_0} \rangle &= -(p^* - 1) \int_{B_\delta(x_0)} f_{\epsilon,x_0} H_{\epsilon,x_0} \left[\langle x - x_0, \frac{\nabla U_{\epsilon,x_0}}{U_{\epsilon,x_0}} \rangle + \frac{N - p}{p} \right] \\ & \quad - \frac{N - p}{p} \int_{B_\delta(x_0)} f_{\epsilon,x_0} G_{\epsilon,x_0} + o_\epsilon(1) \end{aligned} \tag{3.10}$$

as $\epsilon \rightarrow 0$ in view of (2.8)-(2.9) and (2.11). Since there holds

$$\begin{aligned} & \frac{p(p-1)}{N-p} C_0^{1-p} C_1^{-\frac{p^2}{N-p}} \int_{B_\delta(x_0)} f_{\epsilon,x_0} H_{\epsilon,x_0} \left[\langle x - x_0, \frac{\nabla U_{\epsilon,x_0}}{U_{\epsilon,x_0}} \rangle + \frac{N - p}{p} \right] \\ &= \epsilon^p \int_{B_\delta(x_0)} H_{\epsilon,x_0} \frac{(p-1)\epsilon^p - |x - x_0|^{\frac{p}{p-1}}}{(\epsilon^p + |x - x_0|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} \\ &= \int_{B_{\frac{\delta}{\epsilon^{p-1}}}(0)} H_{\epsilon,x_0} (\epsilon^{p-1}y + x_0) \frac{(p-1) - |y|^{\frac{p}{p-1}}}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} \\ &\rightarrow \int_{\mathbb{R}^N} \frac{(p-1) - |y|^{\frac{p}{p-1}}}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} H_\lambda(x_0, x_0) \end{aligned}$$

as $\epsilon \rightarrow 0$ in view of (2.4), (2.11) and the Lebesgue convergence Theorem, we can insert (3.10) into (3.9) and as $\epsilon \rightarrow 0$ get the validity of

$$C_0 H_\lambda(x_0, x_0) = \int_{B_\delta(x_0)} \lambda G_\lambda(x, x_0)^p dx$$

$$\begin{aligned}
 &+ \int_{\partial B_\delta(x_0)} \left(\frac{\delta}{p} |\nabla G_\lambda(x, x_0)|^p - \delta |\nabla G_\lambda(x, x_0)|^{p-2} (\partial_\nu G_\lambda(x, x_0))^2 \right. \\
 &\quad \left. - \frac{\lambda \delta}{p} G_\lambda^p(x, x_0) - \frac{N-p}{p} G_\lambda(x, x_0) |\nabla G_\lambda(x, x_0)|^{p-2} \partial_\nu G_\lambda(x, x_0) \right) d\sigma(x)
 \end{aligned}$$

in view of (2.20) and $\lim_{\epsilon \rightarrow 0} G_{\epsilon, x_0} = G_\lambda(\cdot, x_0)$ in $L^p(\Omega)$ if $\lambda \neq 0$, as it follows by (2.18) and $\bar{q}^* > p$ thanks to $N < 2p \leq p^2$, where

$$C_0 = (p^* - 1) \frac{N-p}{p(p-1)} C_0^{p-1} C_1^{\frac{p^2}{N-p}} \int_{\mathbb{R}^N} \frac{|y|^{\frac{p}{p-1}} - (p-1)}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}}.$$

Concerning the sign of the constant C_0 , observe that

$$\begin{aligned}
 \int_{\mathbb{R}^N} \frac{|y|^{\frac{p}{p-1}}}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} &= -\frac{p-1}{pN + p - N} \int_{\mathbb{R}^N} \langle y, \nabla(1 + |y|^{\frac{p}{p-1}})^{\frac{N}{p}-N-1} \rangle \\
 &= \frac{N(p-1)}{pN + p - N} \int_{\mathbb{R}^N} (1 + |y|^{\frac{p}{p-1}})^{\frac{N}{p}-N-1}
 \end{aligned}$$

and then

$$\int_{\mathbb{R}^N} \frac{|y|^{\frac{p}{p-1}}}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} = \frac{N(p-1)}{p} \int_{\mathbb{R}^N} \frac{1}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}},$$

which implies $C_0 > 0$ in view of

$$\int_{\mathbb{R}^N} \frac{|y|^{\frac{p}{p-1}} - (p-1)}{(1 + |y|^{\frac{p}{p-1}})^{N+2-\frac{N}{p}}} = \frac{(N-p)(p-1)}{p} \int_{\mathbb{R}^N} (1 + |y|^{\frac{p}{p-1}})^{\frac{N}{p}-N-2} > 0.$$

The proof of (3.8) is complete. \square

4. The blow-up approach

Following [9] let us introduce the following blow-up procedure. Letting $\lambda_n = \lambda_* + \frac{1}{n}$, we have that $S_{\lambda_n} < S_0 = S_{\lambda_*}$ and then S_{λ_n} is achieved by a nonnegative $u_n \in W_0^{1,p}(\Omega)$ which, up to a normalization, satisfies

$$-\Delta_p u_n = \lambda_n u_n^{p-1} + u_n^{p^*-1} \text{ in } \Omega, \quad \int_{\Omega} u_n^{p^*} = S_{\lambda_n}^{\frac{N}{p}}. \tag{4.1}$$

Since $\lambda_* < \lambda_1$, by (4.1) the sequence u_n is uniformly bounded in $W_0^{1,p}(\Omega)$ and then, up to a subsequence, $u_n \rightharpoonup u_0 \geq 0$ in $W_0^{1,p}(\Omega)$ and a.e. in Ω as $n \rightarrow +\infty$. Since

$$Q_{\lambda_n}(u) = Q_{\lambda_*}(u) - \frac{1}{n} \frac{\|u_n\|_p^p}{\|u_n\|_{p^*}^p} \geq S_0 - \frac{C}{n}$$

for some $C > 0$ thanks to the Hölder’s inequality, we deduce that

$$\lim_{n \rightarrow +\infty} S_{\lambda_n} = S_0. \tag{4.2}$$

By letting $n \rightarrow +\infty$ in (4.1) we deduce that $u_0 \in W_0^{1,p}(\Omega)$ solves

$$-\Delta_p u_0 = \lambda_* u_0^{p-1} + u_0^{p^*-1} \text{ in } \Omega, \quad \int_{\Omega} u_0^{p^*} \leq S_0^{\frac{N}{p}},$$

thanks to $u_n \rightarrow u_0$ a.e. in Ω as $n \rightarrow +\infty$ and the Fatou convergence Theorem, and then

$$S_0 \leq Q_{\lambda_*}(u_0) = \left(\int_{\Omega} u_0^{p^*} \right)^{\frac{p}{N}} \leq S_0$$

if $u_0 \neq 0$. Since $S_{\lambda_*} = S_0$ would be achieved by u_0 if $u_0 \neq 0$, assumption (1.2) is crucial to guarantee $u_0 = 0$ and then

$$u_n \rightharpoonup 0 \text{ in } W_0^{1,p}(\Omega), \quad u_n \rightarrow 0 \text{ in } L^q(\Omega) \text{ for } 1 \leq q < p^* \text{ and a.e. in } \Omega \tag{4.3}$$

in view of the Sobolev embedding Theorem. Since by elliptic regularity theory [6,16,22,23] and the strong maximum principle [24] $0 < u_n \in C^{1,\alpha}(\bar{\Omega})$ for some $\alpha \in (0, 1)$, we can start a blow-up approach to describe the behavior of u_n since $\|u_n\|_{\infty} \rightarrow +\infty$ as $n \rightarrow +\infty$, as it follows by (4.3) and $\int_{\Omega} u_n^{p^*} = S_{\lambda_n}^{\frac{N}{p}} \rightarrow S_0^{\frac{N}{p}}$ as $n \rightarrow +\infty$.

Letting $x_n \in \Omega$ so that $u_n(x_n) = \max_{\Omega} u_n$, define the blow-up speed as $\mu_n = [u_n(x_n)]^{-\frac{p}{N-p}} \rightarrow 0$ as $n \rightarrow +\infty$ and the blow-up profile

$$U_n(y) = \mu_n^{\frac{N-p}{p}} u_n(\mu_n y + x_n), \quad y \in \Omega_n = \frac{\Omega - x_n}{\mu_n}, \tag{4.4}$$

which satisfies

$$-\Delta_p U_n = \lambda_n \mu_n^p U_n^{p-1} + U_n^{p^*-1} \text{ in } \Omega_n, \quad U_n = 0 \text{ on } \partial\Omega_n \tag{4.5}$$

with $0 < U_n \leq U_n(0) = 1$ in Ω_n and

$$\sup_{n \in \mathbb{N}} \left[\int_{\Omega_n} |\nabla U_n|^p + \int_{\Omega_n} U_n^{p^*} \right] < +\infty.$$

Since U_n is uniformly bounded in $C^{1,\alpha}(A \cap \Omega_n)$ for all $A \subset \subset \mathbb{R}^N$ by elliptic estimates [6,16,22,23], we get that, up to a subsequence, $U_n \rightarrow U$ in $C_{loc}^1(\bar{\Omega}_{\infty})$, where Ω_{∞} is an

half-space with $\text{dist}(0, \partial\Omega_\infty) = L \in (0, \infty]$ in view of $1 = U_n(0) - U_n(y) \leq C|y|$ for $y \in B_2(0) \cap \partial\Omega_n$ and $U \in D^{1,p}(\Omega_\infty)$ solves

$$-\Delta_p U = U^{p^*-1} \text{ in } \Omega_\infty, \quad U = 0 \text{ on } \partial\Omega_\infty, \quad 0 < U \leq U(0) = 1 \text{ in } \Omega_\infty.$$

Since $L < +\infty$ would provide $U \in D_0^{1,p}(\Omega_\infty)$, by [17] one would get $U = 0$, in contradiction with $U(0) = 1$. Since

$$\lim_{n \rightarrow +\infty} \frac{\text{dist}(x_n, \partial\Omega)}{\mu_n} = \lim_{n \rightarrow +\infty} \text{dist}(0, \partial\Omega_n) = +\infty, \tag{4.6}$$

by [5,21,25] we have that U coincides with $U_\infty = (1 + \Lambda|y|^{\frac{p}{p-1}})^{-\frac{N-p}{p}}$, $\Lambda = C_1^{-\frac{p^2}{(N-p)(p-1)}}$ (by (2.5) with $x_0 = 0$ and $\epsilon = C_1^{\frac{p}{(N-p)(p-1)}}$ to have $U_\infty(0) = 1$). Since

$$U_n(y) = \mu_n^{\frac{N-p}{p}} u_n(\mu_n y + x_n) \rightarrow (1 + \Lambda|y|^{\frac{p}{p-1}})^{-\frac{N-p}{p}} \text{ uniformly in } B_R(0) \tag{4.7}$$

as $n \rightarrow +\infty$ for all $R > 0$, in particular there holds

$$\lim_{R \rightarrow +\infty} \lim_{n \rightarrow +\infty} \int_{B_{R\mu_n}(x_n)} u_n^{p^*} = \int_{\mathbb{R}^N} U_\infty^{p^*} = S_0^{\frac{N}{p}}. \tag{4.8}$$

Contained in (4.1)-(4.2), the energy information $\lim_{n \rightarrow +\infty} \int_{\Omega} u_n^{p^*} = S_0^{\frac{N}{p}}$ combines with (4.8) to give

$$\lim_{R \rightarrow +\infty} \lim_{n \rightarrow +\infty} \int_{\Omega \setminus B_{R\mu_n}(x_n)} u_n^{p^*} = 0, \tag{4.9}$$

a property which will simplify the blow-up description of u_n . Up to a subsequence, let us assume $x_n \rightarrow x_0 \in \bar{\Omega}$ as $n \rightarrow +\infty$.

The proof of the implication (iii) \Rightarrow (i) in Theorem 1.1 proceeds through the 5 steps that will be developed below. The main technical point is to establish a comparison between u_n and the bubble

$$U_n(x) = \frac{\mu_n^{\frac{N-p}{p(p-1)}}}{(\mu_n^{\frac{p}{p-1}} + \Lambda|x - x_n|^{\frac{p}{p-1}})^{\frac{N-p}{p}}}$$

in the form $u_n \leq CU_n$ in Ω , no matter x_n tends to $\partial\Omega$ or not. Thanks to such a fundamental estimate, we will first apply some Pohozaev identity in the whole Ω_n to exclude the boundary blow-up $d_n = \text{dist}(x_n, \partial\Omega) \rightarrow 0$ as $n \rightarrow +\infty$. In the interior case, still by a Pohozaev identity on $B_\delta(x_n)$ as $n \rightarrow +\infty$ and $\delta \rightarrow 0$, we will obtain an

information on the limiting blow-up point $x_0 = \lim_{n \rightarrow +\infty} x_n \in \Omega$ in the form $H_{\lambda_*}(x_0, x_0) = 0$ and then the property $H_\lambda(x_0, x_0) > H_{\lambda_*}(x_0, x_0) = 0$ for $\lambda > \lambda_*$ will follow by the monotonicity of $H_\lambda(x_0, x_0)$.

Step 1. There holds $u_n \rightarrow 0$ in $C_{loc}(\bar{\Omega} \setminus \{x_0\})$ as $n \rightarrow +\infty$, where $x_0 = \lim_{n \rightarrow +\infty} x_n \in \bar{\Omega}$.

First observe that

$$u_n \rightarrow 0 \text{ in } L_{loc}^{p^*}(\bar{\Omega} \setminus \{x_0\}) \tag{4.10}$$

as $n \rightarrow +\infty$ in view of (4.9) and we are then concerned with establishing the uniform convergence by a Moser iterative argument. Given a compact set $K \subset \bar{\Omega} \setminus \{x_0\}$, consider $\eta \in C_0^\infty(\mathbb{R}^N \setminus \{x_0\})$ be a cut-off function with $0 \leq \eta \leq 1$ and $\eta = 1$ in K . Since $u_n = 0$ on $\partial\Omega$, use $\eta^p u_n^\beta$, $\beta \geq 1$, as a test function in (4.1) to get

$$\begin{aligned} \frac{\beta p^p}{(\beta - 1 + p)^p} \int_{\Omega} \eta^p |\nabla w_n|^p &\leq \frac{p^p}{(\beta - 1 + p)^{p-1}} \int_{\Omega} \eta^{p-1} |\nabla \eta| w_n |\nabla w_n|^{p-1} \\ &\quad + \int_{\Omega} \lambda_n \eta^p w_n^p + \int_{\Omega} \eta^p u_n^{p^* - p} w_n^p \end{aligned}$$

in terms of $w_n = u_n^{\frac{\beta-1+p}{p}}$ and then by the Young inequality

$$\int_{\Omega} \eta^p |\nabla w_n|^p \leq C \beta^p \left(\int_{\Omega} |\nabla \eta|^p w_n^p + \int_{\Omega} \eta^p w_n^p + \int_{\Omega} \eta^p u_n^{p^* - p} w_n^p \right) \tag{4.11}$$

for some $C > 0$. Since by the Hölder inequality

$$\int_{\Omega} \eta^p u_n^{p^* - p} w_n^p \leq C \left(\int_{\Omega \cap \text{supp } \eta} u_n^{p^*} \right)^{\frac{p}{N}} \|\eta w_n\|_{p^*}^p = o(\|\eta w_n\|_{p^*}^p)$$

as $n \rightarrow +\infty$ in view of (4.10) and $\Omega \cap \text{supp } \eta \subset \subset \bar{\Omega} \setminus \{x_0\}$, by (4.11) and the Sobolev embedding Theorem we deduce that

$$\|\eta w_n\|_{p^*}^p \leq C \|w_n\|_p^p = C \int_{\Omega} u_n^{\beta-1+p} \rightarrow 0$$

for all $1 \leq \beta < p^* - p + 1$ in view of (4.3) and then $u_n \rightarrow 0$ in $L^q(K)$ for all $1 \leq q < \frac{Np^*}{N-p}$ as $n \rightarrow +\infty$. We have then established that

$$u_n \rightarrow 0 \text{ in } L_{loc}^q(\bar{\Omega} \setminus \{x_0\}) \tag{4.12}$$

as $n \rightarrow +\infty$ for all $1 \leq q < \frac{Np^*}{N-p}$. Since $\frac{N}{p}(p^* - p) = p^* < \frac{Np^*}{N-p}$, observe that (4.12) now provides that the R.H.S. in the equation (4.1) can be written as $(\lambda_n + u_n^{p^*-p})u_n^{p-1}$ with a bound on the coefficient $\lambda_n + u_n^{p^*-p}$ in $L_{loc}^{q_0}(\bar{\Omega} \setminus \{x_0\})$ for some $q_0 > \frac{N}{p}$. Given compact sets $K \subset \tilde{K} \subset \bar{\Omega} \setminus \{x_0\}$ with $\text{dist}(K, \partial\tilde{K}) > 0$, by [22] we have the estimate $\|u_n\|_{\infty, K} \leq C\|u_n\|_{p, \tilde{K}}$ and then $u_n \rightarrow 0$ in $C(K)$ as $n \rightarrow +\infty$ in view of (4.12) and $p < \frac{Np^*}{N-p}$. The convergence $u_n \rightarrow 0$ in $C_{loc}(\bar{\Omega} \setminus \{x_0\})$ has been then established as $n \rightarrow +\infty$.

Step 2. The following pointwise estimates

$$\lim_{n \rightarrow +\infty} \max_{\Omega} |x - x_n|^{\frac{N-p}{p}} u_n < \infty, \quad \lim_{R \rightarrow +\infty} \lim_{n \rightarrow +\infty} \max_{\Omega \setminus B_{R\mu_n}(x_n)} |x - x_n|^{\frac{N-p}{p}} u_n = 0 \quad (4.13)$$

do hold.

By contradiction and up to a subsequence, assume the existence of $y_n \in \Omega$ such that either

$$|x_n - y_n|^{\frac{N-p}{p}} u_n(y_n) = \max_{\Omega} |x - x_n|^{\frac{N-p}{p}} u_n \rightarrow +\infty \quad (4.14)$$

as $n \rightarrow +\infty$ or

$$\max_{\Omega} |x - x_n|^{\frac{N-p}{p}} u_n \leq C_0, \quad |x_n - y_n|^{\frac{N-p}{p}} u_n(y_n) = \max_{\Omega \setminus B_{R_n\mu_n}(x_n)} |x - x_n|^{\frac{N-p}{p}} u_n \geq \delta > 0 \quad (4.15)$$

for some $R_n \rightarrow +\infty$ as $n \rightarrow +\infty$. Setting $\nu_n = [u_n(y_n)]^{-\frac{p}{N-p}}$, there hold $\frac{|x_n - y_n|}{\nu_n} \rightarrow +\infty$ in case (4.14), $\frac{|x_n - y_n|}{\nu_n} \in [\delta^{\frac{p}{N-p}}, C_0^{\frac{p}{N-p}}]$ in case (4.15) and $\nu_n \rightarrow 0$ as $n \rightarrow +\infty$, since $x_n - y_n \rightarrow 0$ as $n \rightarrow +\infty$ when (4.15) holds thanks to Step 1. Up to a further subsequence, let us assume that $\frac{x_n - y_n}{\nu_n} \rightarrow p$ as $n \rightarrow +\infty$, where $p = +\infty$ in case (4.14) and $p \in \mathbb{R}^N \setminus \{0\}$ in case (4.15). Since $(\frac{|x_n - y_n|}{\mu_n})^{\frac{N-p}{p}} \geq (\frac{|x_n - y_n|}{\mu_n})^{\frac{N-p}{p}} U_n(\frac{y_n - x_n}{\mu_n}) = |x_n - y_n|^{\frac{N-p}{p}} u_n(y_n)$ in view of (4.5), where U_n is given by (4.4), then $\frac{|x_n - y_n|}{\mu_n} \rightarrow +\infty$ as $n \rightarrow +\infty$ also in case (4.14). Setting $V_n(y) = \nu_n^{\frac{N-p}{p}} u_n(\nu_n y + y_n)$ for $y \in \tilde{\Omega}_n = \frac{\Omega - y_n}{\nu_n}$, then $V_n(0) = 1$ and in $\tilde{\Omega}_n$ there hold:

$$\begin{aligned} V_n(y) &\leq \nu_n^{\frac{N-p}{p}} |\nu_n y + y_n - x_n|^{-\frac{N-p}{p}} |x_n - y_n|^{\frac{N-p}{p}} u_n(y_n) = \left(\frac{|x_n - y_n|}{|\nu_n y + y_n - x_n|}\right)^{\frac{N-p}{p}} \\ &\leq 2^{\frac{N-p}{p}} \end{aligned} \quad (4.16)$$

for $|y| \leq \frac{1}{2} \frac{|x_n - y_n|}{\nu_n}$ in case (4.14) and

$$\left|y - \frac{x_n - y_n}{\nu_n}\right|^{\frac{N-p}{p}} V_n(y) = |\nu_n y + y_n - x_n|^{\frac{N-p}{p}} u_n(\nu_n y + y_n) \leq C_0 \quad (4.17)$$

in case (4.15). Since

$$-\Delta_p V_n = \lambda_n \nu_n^p V_n^{p-1} + V_n^{p^*-1} \text{ in } \tilde{\Omega}_n, \quad V_n = 0 \text{ on } \partial\tilde{\Omega}_n,$$

by (4.16)-(4.17) and standard elliptic estimates [6,16,22,23] we get that V_n is uniformly bounded in $C^{1,\alpha}(A \cap \tilde{\Omega}_n)$ for all $A \subset\subset \mathbb{R}^N \setminus \{p\}$. Up to a subsequence, we have that $V_n \rightarrow V$ in $C^1_{loc}(\tilde{\Omega}_\infty \setminus \{p\})$, where Ω_∞ is an half-space with $\text{dist}(0, \partial\Omega_\infty) = L$. Since $p \neq 0$, there hold $B_{\frac{|p|}{2}}(0) \subset\subset \mathbb{R}^N \setminus \{p\}$ and $1 = V_n(0) - V_n(y) \leq C|y|$ for $y \in B_{\frac{|p|}{2}}(0) \cap \partial\tilde{\Omega}_n$, leading to $L \in (0, \infty]$. Since $V \geq 0$ solves $-\Delta_p V = V^{p^*-1}$ in Ω_∞ , by the strong maximum principle [24] we deduce that $V > 0$ in Ω_∞ in view of $V(0) = 1$ thanks to $0 \in \Omega_\infty$. Setting $M = \min\{L, |p|\}$, by $\frac{|x_n - y_n|}{\mu_n} \rightarrow +\infty$ as $n \rightarrow +\infty$ we have that $B_{\frac{M}{2}\nu_n}(y_n) \subset \Omega \setminus B_{R\mu_n}(x_n)$ for all $R > 0$ provided n is sufficiently large (depending on R) and then

$$\int_{\Omega \setminus B_{R\mu_n}(x_n)} u_n^{p^*} \geq \int_{B_{\frac{M}{2}\nu_n}(y_n)} u_n^{p^*} = \int_{B_{\frac{M}{2}}(0)} V_n^{p^*} \rightarrow \int_{B_{\frac{M}{2}}(0)} V^{p^*} > 0,$$

in contradiction with (4.9). The proof of (4.13) is complete.

Step 3. There exists $C > 0$ so that

$$u_n \leq \frac{C\mu_n^{\frac{N-p}{p(p-1)}}}{(\mu_n^{\frac{p}{p-1}} + \Lambda|x - x_n|^{\frac{p}{p-1}})^{\frac{N-p}{p}}} \text{ in } \Omega \tag{4.18}$$

does hold for all $n \in \mathbb{N}$.

Since (4.18) does already hold in $B_{R\mu_n}(x_n)$ for all $R > 0$ thanks to (4.7), notice that (4.18) is equivalent to establish the estimate

$$u_n \leq \frac{C\mu_n^{\frac{N-p}{p(p-1)}}}{|x - x_n|^{\frac{N-p}{p-1}}} \text{ in } \Omega \setminus B_{R\mu_n}(x_n) \tag{4.19}$$

for some $C, R > 0$ and all $n \in N$. Let us first prove the following weaker form of (4.19): given $0 < \eta < \frac{N-p}{p(p-1)}$ there exist $C, R > 0$ so that

$$u_n \leq \frac{C\mu_n^{\frac{N-p}{p(p-1)} - \eta}}{|x - x_n|^{\frac{N-p}{p-1} - \eta}} \text{ in } \Omega \setminus B_{R\mu_n}(x_n) \tag{4.20}$$

does hold for all $n \in N$. Since $|x - x_n|^{\eta - \frac{N-p}{p-1}}$ satisfies

$$-\Delta_p |x - x_n|^{\eta - \frac{N-p}{p-1}} = \eta(p-1) \left(\frac{N-p}{p-1} - \eta\right)^{p-1} |x - x_n|^{\eta(p-1) - N},$$

we have that $\Phi_n = C \frac{\mu_n^{\frac{N-p}{p(p-1)} - \eta} + M_n}{|x - x_n|^{\frac{N-p}{p-1} - \eta}}$, where $\rho, C > 0$ and $M_n = \sup_{\Omega \cap \partial B_\rho(x_0)} u_n$, satisfies

$$\begin{aligned}
 & -\Delta_p \Phi_n - \left(\lambda_n + \frac{\delta}{|x - x_n|^p}\right) \Phi_n^{p-1} \\
 & = \left[\eta(p-1) \left(\frac{N-p}{p-1} - \eta\right)^{p-1} - (\lambda_n |x - x_n|^p + \delta) \right] \frac{\Phi_n^{p-1}}{|x - x_n|^p} \\
 & \geq 0 \quad \text{in } \Omega \cap B_\rho(x_0) \setminus \{x_n\}
 \end{aligned}$$

provided ρ and δ are sufficiently small (depending on η). Taking $R > 0$ large so that $u_n^{p^* - p} \leq \frac{\delta}{|x - x_n|^p}$ in $\Omega \setminus B_{R\mu_n}(x_n)$ for all n large thanks to (4.13), we have that

$$-\Delta_p u_n - \left(\lambda_n + \frac{\delta}{|x - x_n|^p}\right) u_n^{p-1} = (u_n^{p^* - p} - \frac{\delta}{|x - x_n|^p}) u_n^{p-1} \leq 0 \quad \text{in } \Omega \setminus B_{R\mu_n}(x_n).$$

By (4.7) on $\partial B_{R\mu_n}(x_n)$ it is easily seen that $u_n \leq \Phi_n$ on the boundary of $\Omega \cap B_\rho(x_0) \setminus B_{R\mu_n}(x_n)$ for some $C > 0$, and then by Proposition 2.3 one deduces the validity of

$$u_n \leq C \frac{\mu_n^{\frac{N-p}{p(p-1)} - \eta} + M_n}{|x - x_n|^{\frac{N-p}{p-1} - \eta}} \tag{4.21}$$

in $\Omega \cap B_\rho(x_0) \setminus B_{R\mu_n}(x_n)$. Setting $A = \Omega \setminus B_\rho(x_0)$, observe that the function $v_n = \frac{u_n}{M_n}$ satisfies

$$-\Delta_p v_n - \lambda_n v_n^{p-1} = f_n \text{ in } \Omega, \quad v_n = 0 \text{ on } \partial\Omega, \quad \sup_{\Omega \cap \partial B_\rho(x_0)} v_n = 1, \tag{4.22}$$

where $f_n = \frac{u_n^{p^* - 1}}{M_n^{p-1}} = u_n^{\frac{p^2}{N-p}} v_n^{p-1}$. Letting g_n be the p -harmonic function in A so that $g_n = v_n$ on ∂A , observe that $\|g_n\|_\infty = 1$ in view of $0 \leq v_n \leq 1$ on ∂A . Since by Step 1 there holds

$$a_n = \lambda_n + u_n^{\frac{p^2}{N-p}} \rightarrow \lambda_* \quad \text{in } L^\infty(A)$$

as $n \rightarrow +\infty$ with $\lambda_* < \lambda_1(\Omega) < \lambda_1(A)$, by Proposition 2.2 we deduce that $\sup_{n \in \mathbb{N}} \|v_n\|_{p-1, A} < +\infty$ and then $\sup_{n \in \mathbb{N}} \|f_n\|_{1, A} < +\infty$ in view of Step 1. Letting w_n the solution of

$$-\Delta_p w_n = f_n \text{ in } A, \quad w_n = 0 \text{ on } \partial A,$$

by Proposition 2.2 we also deduce that $\sup_{n \in \mathbb{N}} \|v_n - w_n\|_{\infty, A} < +\infty$ thanks to $N < 2p$. Since by the Sobolev embedding Theorem $\sup_{n \in \mathbb{N}} \|w_n\|_{q, A} < +\infty$ for all $1 \leq q < \bar{q}^*$ in view of Proposition 2.2 and $\sup_{n \in \mathbb{N}} \|f_n\|_{1, A} < +\infty$, similar estimates hold for v_n and then $\sup_{n \in \mathbb{N}} \|f_n\|_{q_0, A} < +\infty$ for some $q_0 > \frac{N}{p}$ in view of $N < 2p$. By elliptic estimates [22] we get that $\sup_{n \in \mathbb{N}} \|w_n\|_{\infty, A} < +\infty$ and in turn $\sup_{n \in \mathbb{N}} \|v_n\|_{\infty, A} < +\infty$, or equivalently

$$\sup_{\Omega \setminus B_\rho(x_0)} u_n \leq C \sup_{\Omega \cap \partial B_\rho(x_0)} u_n \tag{4.23}$$

for some $C > 0$. Thanks to (4.23) one can extend the validity of (4.21) from $\Omega \cap B_\rho(x_0) \setminus B_{R\mu_n}(x_n)$ to $\Omega \setminus B_{R\mu_n}(x_n)$. In order to establish (4.20), we claim that M_n in (4.21) satisfies

$$M_n = o(\mu_n^{\frac{N-p}{p(p-1)}-\eta}) \tag{4.24}$$

for all $0 < \eta < \frac{N-p}{p(p-1)}$.

Indeed, by contradiction assume that there exist $0 < \bar{\eta} < \frac{N-p}{p(p-1)}$ and a subsequence so that

$$\mu_n^{\frac{N-p}{p(p-1)}-\bar{\eta}} \leq CM_n \tag{4.25}$$

for some $C > 0$. Since $v_n = O(|x - x_n|^{-\frac{N-p}{p-1}+\bar{\eta}})$ uniformly in $\Omega \setminus B_{R\mu_n}(x_n)$ in view of (4.21) and (4.25), we have that v_n and then $f_n = u_n^{\frac{p^2}{N-p}} v_n^{p-1}$ are uniformly bounded in $C_{loc}(\bar{\Omega} \setminus \{x_0\})$ and by elliptic estimates [6,16,22,23] $v_n \rightarrow v$ in $C_{loc}^1(\bar{\Omega} \setminus \{x_0\})$ as $n \rightarrow +\infty$, up to a further subsequence, where $v \neq 0$ in view of $\sup_{\Omega \cap \partial B_\rho(x_0)} v = \lim_{n \rightarrow +\infty} \sup_{\Omega \cap \partial B_\rho(x_0)} v_n = 1$.

Moreover, notice that $\lim_{n \rightarrow +\infty} \|f_n\|_1 = 0$ would imply $v_n \rightarrow v$ in $W_0^{1,q}(\Omega)$ for all $1 \leq q < \bar{q}$ and in $L^s(\Omega)$ for all $1 \leq s < \bar{q}^*$ as $n \rightarrow +\infty$ in view of Proposition 2.2, where v is a solution of

$$-\Delta_p v - \lambda_* v^{p-1} = 0 \quad \text{in } \Omega. \tag{4.26}$$

Letting

$$T_l(s) = \begin{cases} |s| & \text{if } |s| \leq l \\ \pm l & \text{if } \pm s > l \end{cases}$$

and using $T_l(v_n) \in W_0^{1,p}(\Omega)$ as a test function in (4.22), one would get

$$\int_{\{|v_n| \leq l\}} |\nabla v_n|^p \leq \lambda_n \int_{\Omega} v_n^p + l \|f_n\|_1 \rightarrow \lambda_* \int_{\Omega} v^p$$

as $n \rightarrow +\infty$ in view of $\bar{q}^* > p$ and then deduce

$$\int_{\Omega} |\nabla v|^p \leq \lambda_* \int_{\Omega} v^p < +\infty$$

as $l \rightarrow +\infty$. Since $v \in W_0^{1,p}(\Omega)$ solves (4.26) with $\lambda_* < \lambda_1$, one would have $v = 0$, in contradiction with $\sup_{\Omega \cap \partial B_\rho(x_0)} v = 1$. Once

$$\liminf_{n \rightarrow +\infty} \|f_n\|_1 > 0 \tag{4.27}$$

has been established, by (4.4), (4.7) and (4.21) observe that

$$\int_{B_{R\mu_n}(x_n)} f_n = \int_{B_{R\mu_n}(x_n)} \frac{u_n^{p^*-1}}{M_n^{p-1}} = \frac{\mu_n^{\frac{N-p}{p}}}{M_n^{p-1}} \int_{B_R(0)} U_n^{p^*-1} = O\left(\frac{\mu_n^{\frac{N-p}{p}}}{M_n^{p-1}}\right) \tag{4.28}$$

and

$$\int_{\Omega \setminus B_{R\mu_n}(x_n)} f_n = L_n^{\frac{p^2}{N-p}} \left(\frac{L_n}{M_n}\right)^{p-1} O\left(\mu_n^{(p^*-1)\eta - \frac{p}{p-1}} \log \frac{1}{\mu_n} + 1\right) \tag{4.29}$$

where $L_n = \mu_n^{\frac{N-p}{p(p-1)} - \eta} + M_n$. Setting $\eta_0 = \frac{p}{(p-1)(p^*-1)}$, then (4.24) necessarily holds for $\eta \in (\eta_0, \frac{N-p}{p(p-1)})$ since otherwise $L_n = O(M_n)$ and (4.28)-(4.29) would provide $\lim_{n \rightarrow +\infty} \|f_n\| = 0$ along a subsequence thanks to $\lim_{n \rightarrow +\infty} M_n = 0$, in contradiction with (4.27). Notice that (4.24) holds for $\eta = \eta_0$ too, since otherwise the conclusion $\lim_{n \rightarrow +\infty} \|f_n\| = 0$ would follow as above thanks to $L_n = O(\mu_n^{\frac{N-p}{p(p-1)} - \eta})$ for $\eta \in (\eta_0, \frac{N-p}{p(p-1)})$. Setting $\eta_k = (\frac{p^2}{N-p-N+p})^k \eta_0$, arguing as above (4.24) can be established for $\eta \in [\eta_{k+1}, \eta_k]$, $k \geq 0$, by using the validity of (4.24) for $\eta \in [\eta_k, \frac{N-p}{p-1})$ in view of the relation

$$(p^* - 1)\eta_{k+1} - \frac{p}{p-1} + \frac{p^2}{N-p} \left[\frac{N-p}{p(p-1)} - \eta_k \right] = 0.$$

Since $\frac{p^2}{N-p-N+p} < 1$ for $p < N$, we have that $\eta_k \rightarrow 0$ as $k \rightarrow +\infty$ and then (4.24) is proved for all $0 < \eta < \frac{N-p}{p(p-1)}$, in contradiction with (4.25). Therefore, we have established (4.24) and the validity of (4.20) follows.

In order to establish (4.19), let us repeat the previous argument for $v_n = \mu_n^{-\frac{N-p}{p(p-1)}} u_n$, where v_n solves

$$-\Delta_p v_n - \lambda_n v_n^{p-1} = f_n \text{ in } \Omega, \quad v_n = 0 \text{ on } \partial\Omega, \tag{4.30}$$

with $f_n = \mu_n^{-\frac{N-p}{p}} u_n^{p^*-1}$. Notice that f_n satisfies

$$f_n \leq \frac{C_0 \mu_n^{\frac{p}{p-1} - (p^*-1)\eta}}{|x - x_n|^{N + \frac{p}{p-1} - (p^*-1)\eta}} \text{ in } \Omega \setminus B_{R\mu_n}(x_n) \tag{4.31}$$

for some $C_0 > 0$ in view of (4.20) and then, by arguing as in (4.28),

$$\int_{\Omega} f_n = O(1) + O\left(\int_{\Omega \setminus B_{R\mu_n}(x_n)} \frac{\mu_n^{\frac{p}{p-1} - (p^*-1)\eta}}{|x - x_n|^{N + \frac{p}{p-1} - (p^*-1)\eta}}\right) = O(1) \tag{4.32}$$

for $0 < \eta < \frac{p}{(p-1)(p^*-1)} = \frac{p(N-p)}{(p-1)(Np-N+p)}$. Letting h_n be the solution of

$$-\Delta_p h_n = f_n \text{ in } \Omega, \quad h_n = 0 \text{ on } \partial\Omega,$$

by (4.32) and Proposition 2.2 we deduce that $\sup_{n \in \mathbb{N}} \|v_n - h_n\|_\infty < +\infty$ thanks to $N < 2p$, or equivalently

$$\|u_n - \mu_n^{\frac{N-p}{p(p-1)}} h_n\|_\infty = O(\mu_n^{\frac{N-p}{p(p-1)}}). \tag{4.33}$$

For $\alpha > N$ the radial function

$$W(y) = (\alpha - N)^{-\frac{1}{p-1}} \int_{|y|}^\infty \frac{(t^{\alpha-N} - 1)^{\frac{1}{p-1}}}{t^{\frac{\alpha-1}{p-1}}} dt$$

is a positive and strictly decreasing solution of $-\Delta_p W = |y|^{-\alpha}$ in $\mathbb{R}^N \setminus B_1(0)$ so that

$$\lim_{|y| \rightarrow \infty} |y|^{\frac{N-p}{p-1}} W(y) = \frac{p-1}{N-p} (\alpha - N)^{-\frac{1}{p-1}} > 0. \tag{4.34}$$

Taking $0 < \eta < \frac{p}{(p-1)(p^*-1)}$ to ensure $\alpha := N + \frac{p}{p-1} - (p^* - 1)\eta > N$, then $w_n(x) = \mu_n^{-\frac{N-p}{p-1}} W(\frac{x-x_n}{\mu_n})$ satisfies

$$-\Delta_p w_n = \frac{\mu_n^{\frac{p}{p-1} - (p^* - 1)\eta}}{|x - x_n|^{N + \frac{p}{p-1} - (p^* - 1)\eta}} \text{ in } \mathbb{R}^N \setminus B_1(x_n).$$

Since

$$h_n(x) = \mu_n^{-\frac{N-p}{p(p-1)}} u_n(x) + O(1) = \mu_n^{-\frac{N-p}{p-1}} U_n(\frac{x - x_n}{\mu_n}) + O(1) \leq C_1 w_n(x)$$

for some $C_1 > 0$ and for all $x \in \partial B_{R\mu_n}(x_n)$ in view of (4.7), (4.33) and $W(R) > 0$, we have that $\Phi_n = Cw_n$ satisfies

$$-\Delta_p \Phi_n \geq f_n \text{ in } \Omega \setminus B_{R\mu_n}(x_n), \quad \Phi_n \geq h_n \text{ on } \partial\Omega \cup \partial B_{R\mu_n}(x_n)$$

for $C = C_0^{\frac{1}{p-1}} + C_1$ thanks to (4.31), and then by weak comparison principle we deduce that

$$h_n \leq \Phi_n \leq \frac{C}{|x - x_n|^{\frac{N-p}{p-1}}} \text{ in } \Omega \setminus B_{R\mu_n}(x_n) \tag{4.35}$$

for some $C > 0$ in view of (4.34). Inserting (4.35) into (4.33) we finally deduce the validity of (4.19)

Step 4. There holds $x_0 \notin \partial\Omega$.

Assume by contradiction $x_0 \in \partial\Omega$ and set $\hat{x} = x_0 - \nu(x_0)$. Let us apply the Pohozaev identity (3.6) to u_n with $c = 1$, $f = 0$ and $x_0 = \hat{x}$ on $D = \Omega$, together with (3.7), to get

$$\int_{\partial\Omega} |\nabla u_n|^p \langle x - \hat{x}, \nu \rangle = \frac{p}{p-1} \lambda_n \int_{\Omega} u_n^p \tag{4.36}$$

in view of $u_n = 0$ and $\nabla u_n = (\partial_\nu u_n)\nu$ on $\partial\Omega$. Since $v_n = \mu_n^{-\frac{N-p}{p(p-1)}} u_n$ solves (4.30) and v_n, f_n are uniformly bounded in $C_{loc}^1(\bar{\Omega} \setminus \{x_0\})$ in view of (4.18) and (4.31) with $\eta = 0$, by elliptic estimates [6,16,22,23] we deduce that v_n is uniformly bounded in $C_{loc}^1(\bar{\Omega} \setminus \{x_0\})$. Fixing $\rho > 0$ small so that $\langle x - \hat{x}, \nu(x) \rangle \geq \frac{1}{2}$ for all $x \in \partial\Omega \cap B_\rho(x_0)$, by (4.18), (4.36) and the C^1 -bound on v_n we have that

$$\int_{\partial\Omega \cap B_\rho(x_0)} |\nabla u_n|^p = O(\lambda_n \int_{\Omega} u_n^p) + \int_{\partial\Omega \setminus B_\rho(x_0)} |\nabla u_n|^p = O(\mu_n^{\frac{N-p}{p-1}}) \tag{4.37}$$

since $\frac{p(N-p)}{p-1} < N$ thanks $N < 2p \leq p^2$. Setting $d_n = \text{dist}(x_n, \partial\Omega)$ and $W_n(y) = d_n^{\frac{N-p}{p}} u_n(d_n y + x_n)$ for $y \in \Omega_n = \frac{\Omega - x_n}{d_n}$, we have that $d_n \rightarrow 0$ and $\Omega_n \rightarrow \Omega_\infty$ as $n \rightarrow +\infty$ where Ω_∞ is an halfspace containing 0 with $\text{dist}(0, \partial\Omega_\infty) = 1$. Setting $\delta_n = \frac{\mu_n}{d_n} \rightarrow 0$ as $n \rightarrow +\infty$ in view of (4.6), the function $G_n = \delta_n^{-\frac{N-p}{p(p-1)}} W_n = \mu_n^{-\frac{N-p}{p(p-1)}} d_n^{\frac{N-p}{p-1}} u_n(d_n y + x_n) \geq 0$ solves

$$-\Delta_p G_n - \lambda_n d_n^p G_n^{p-1} = \tilde{f}_n \text{ in } \Omega_n, \quad G_n = 0 \text{ on } \partial\Omega_n, \tag{4.38}$$

with $\tilde{f}_n = \mu_n^{-\frac{N-p}{p}} d_n^N u_n^{p^*-1}(d_n y + x_n) = d_n^N f_n(d_n y + x_n)$ so that

$$\tilde{f}_n \leq \frac{C \delta_n^{\frac{p}{p-1}}}{|y|^{N+\frac{p}{p-1}}}, \quad G_n \leq \frac{C}{|y|^{\frac{N-p}{p-1}}} \text{ in } \Omega_n \tag{4.39}$$

in view of (4.18) and (4.31) with $\eta = 0$. By (4.39) and elliptic estimates [6,16,22,23] we deduce that $G_n \rightarrow G$ in $C_{loc}^1(\bar{\Omega}_\infty \setminus \{0\})$ as $n \rightarrow +\infty$, where $G \geq 0$ does solve

$$-\Delta_p G = \left(\int_{\mathbb{R}^N} U^{p^*-1} \right) \delta_0 \text{ in } \Omega_\infty, \quad G = 0 \text{ on } \partial\Omega_\infty,$$

in view of (4.38) and

$$\lim_{n \rightarrow +\infty} \int_{B_\epsilon(0)} \tilde{f}_n = \lim_{n \rightarrow +\infty} \int_{B_{\epsilon \frac{d_n}{\mu_n}}(0)} U_n^{p^*-1} = \int_{\mathbb{R}^N} U^{p^*-1} \tag{4.40}$$

for all $\epsilon > 0$ in view of (4.6)-(4.7) and (4.18). By the strong maximum principle [24] we then have that $G > 0$ in Ω_∞ and $\partial_\nu G < 0$ on $\partial\Omega_\infty$. On the other hand, for any $R > 0$ there holds

$$\int_{\partial\Omega_n \cap B_R(0)} |\nabla G_n|^p = \mu_n^{-\frac{N-p}{p-1}} d_n^{\frac{N-1}{p-1}} \int_{\partial\Omega \cap B_{Rd_n}(x_n)} |\nabla u_n|^p = O(d_n^{\frac{N-1}{p-1}})$$

in view of (4.37) and then as $n \rightarrow +\infty$

$$\int_{\partial\Omega_\infty \cap B_R(0)} |\nabla G|^p = 0.$$

We end up with the contradictory conclusion $\nabla G = 0$ on $\partial\Omega_\infty$, and then $x_0 \notin \partial\Omega$.

Step 5. There holds $H_{\lambda_*}(x_0, x_0) = 0$.

Let us apply the Pohozaev identity (3.6) to u_n with $c = 1$ and $f = 0$ on $D = B_\delta(x_0) \subset \Omega$ and (3.7) to get

$$\begin{aligned} \lambda_n \int_{B_\delta(x_0)} u_n^p + \int_{\partial B_\delta(x_0)} \left(\frac{\delta}{p} |\nabla u_n|^p - \delta |\nabla u_n|^{p-2} (\partial_\nu u_n)^2 \right. \\ \left. - \frac{\lambda_n \delta}{p} u_n^p - \frac{N-p}{p} u_n |\nabla u_n|^{p-2} \partial_\nu u_n \right) - \frac{N-p}{Np} \delta \int_{\partial B_\delta(x_0)} u_n^{p^*} = 0. \end{aligned} \tag{4.41}$$

As in the previous Step, up to a subsequence, there holds $G_n = \mu_n^{-\frac{N-p}{p(p-1)}} u_n \rightarrow G$ in $C^1_{loc}(\bar{\Omega} \setminus \{x_0\})$ as $n \rightarrow +\infty$, where $G \geq 0$ satisfies

$$-\Delta_p G - \lambda_* G^{p-1} = \left(\int_{\mathbb{R}^N} U^{p^*-1} \right) \delta_{x_0} \text{ in } \Omega, \quad G = 0 \text{ on } \partial\Omega,$$

as it follows by (4.38) and (4.40) with $d_n = 1$. Arguing as in Proposition 2.4, we can prove that $H = G - \Gamma$ satisfies (2.1) and by Theorem 2.1 it follows that $G = \left(\int_{\mathbb{R}^N} U^{p^*-1} \right)^{\frac{1}{p-1}} G_{\lambda_*}(\cdot, x_0)$. Since $\mu_n^{-\frac{N-p}{p(p-1)}} u_n \rightarrow \left(\int_{\mathbb{R}^N} U^{p^*-1} \right)^{\frac{1}{p-1}} G_{\lambda_*}(\cdot, x_0)$ in $C^1_{loc}(\bar{\Omega} \setminus \{x_0\})$ as $n \rightarrow +\infty$, by letting $n \rightarrow +\infty$ in (4.41) we finally get

$$\begin{aligned} \lambda_* \int_{B_\delta(x_0)} G_{\lambda_*}^p(x, x_0) dx + \int_{\partial B_\delta(x_0)} \left(\frac{\delta}{p} |\nabla G_{\lambda_*}(x, x_0)|^p - \delta |\nabla G_{\lambda_*}(x, x_0)|^{p-2} (\partial_\nu G_{\lambda_*}(x, x_0))^2 \right. \\ \left. - \frac{\lambda_* \delta}{p} G_{\lambda_*}^p(x, x_0) - \frac{N-p}{p} G_{\lambda_*}(x, x_0) |\nabla G_{\lambda_*}(x, x_0)|^{p-2} \partial_\nu G_{\lambda_*}(x, x_0) \right) d\sigma(x) = 0 \end{aligned}$$

and then $H_{\lambda_*}(x_0, x_0) = 0$ by (3.8).

Data availability

No data was used for the research described in the article.

References

- [1] S. Angeloni, P. Esposito, The Green function for p -Laplace operators, preprint, arXiv:2203.01206.
- [2] T. Aubin, Y.Y. Li, On the best Sobolev inequality, *J. Math. Pures Appl.* (9) 78 (4) (1999) 353–387.
- [3] H. Brezis, L. Nirenberg, Positive solutions of nonlinear elliptic equations involving critical Sobolev exponents, *Commun. Pure Appl. Math.* 36 (1983) 437–477.
- [4] L. Damascelli, A. Farina, B. Sciunzi, E. Valdinoci, Liouville results for m -Laplace equations of Lane-Emden-Fowler type, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* 26 (4) (2009) 1099–1119.
- [5] L. Damascelli, S. Merchán, L. Montoro, B. Sciunzi, Radial symmetry and applications for a problem involving the $-\Delta_p(\cdot)$ operator and critical nonlinearity in \mathbb{R}^N , *Adv. Math.* 265 (10) (2014) 313–335.
- [6] E. Dibenedetto, $C^{1+\alpha}$ local regularity of weak solutions of degenerate elliptic equations, *Nonlinear Anal.* 7 (8) (1983) 827–850.
- [7] Z. Djadli, O. Druet, Extremal functions for optimal Sobolev inequalities on compact manifolds, *Calc. Var. Partial Differ. Equ.* 12 (1) (2001) 59–84.
- [8] O. Druet, The best constants problem in Sobolev inequalities, *Math. Ann.* 314 (2) (1999) 327–346.
- [9] O. Druet, Elliptic equations with critical Sobolev exponents in dimension 3, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* 19 (2) (2002) 125–142.
- [10] O. Druet, F. Robert, E. Hebey, *Blow-Up Theory for Elliptic PDEs in Riemannian Geometry*, Mathematical Notes, vol. 45, Princeton University Press, Princeton, NJ, 2004.
- [11] P. Esposito, On some conjectures proposed by Haïm Brezis, *Nonlinear Anal.* 54 (5) (2004) 751–759.
- [12] N. Ghoussoub, F. Robert, Hardy-singular boundary mass and Sobolev-critical variational problems, *Anal. PDE* 10 (5) (2017) 1017–1079.
- [13] N. Ghoussoub, F. Robert, The Hardy-Schrödinger operator with interior singularity: the remaining cases, *Calc. Var. Partial Differ. Equ.* 56 (5) (2017) 149.
- [14] M. Guedda, L. Veron, Quasilinear elliptic equations involving critical Sobolev exponents, *Nonlinear Anal.* 13 (8) (1989) 879–902.
- [15] S. Kichenassamy, L. Veron, Singular solutions of the p -Laplace equation, *Math. Ann.* 275 (1986) 599–616.
- [16] G.M. Lieberman, Boundary regularity for solutions of degenerate elliptic equations, *Nonlinear Anal.* 12 (11) (1988) 1203–1219.
- [17] C. Mercuri, M. Willem, A global compactness result for the p -Laplacian involving critical nonlinearities, *Discrete Contin. Dyn. Syst.* 28 (2) (2010) 469–493.
- [18] R. Schoen, Conformal deformation of a Riemannian metric to constant scalar curvature, *J. Differ. Geom.* 20 (2) (1984) 479–495.
- [19] R. Schoen, S.T. Yau, On the proof of the positive mass conjecture in general relativity, *Commun. Math. Phys.* 65 (1) (1979) 45–76.
- [20] R. Schoen, S.T. Yau, Proof of the positive mass theorem. II, *Commun. Math. Phys.* 79 (2) (1981) 231–260.
- [21] B. Sciunzi, Classification of positive $\mathcal{D}^{1,p}(\mathbb{R}^N)$ -solutions to the critical p -Laplace equation in \mathbb{R}^N , *Adv. Math.* 291 (2016) 12–23.
- [22] J. Serrin, Local behaviour of solutions of quasilinear equations, *Acta Math.* 111 (1964) 247–302.
- [23] P. Tolksdorf, Regularity for a more general class of quasilinear elliptic equations, *J. Differ. Equ.* 51 (1984) 126–150.
- [24] J.L. Vazquez, A strong maximum principle for some quasilinear elliptic equations, *Appl. Math. Optim.* 12 (1984) 191–202.
- [25] J. Vétois, A priori estimates and application to the symmetry of solutions for critical p -Laplace equations, *J. Differ. Equ.* 260 (1) (2016) 149–161.