

A (partial) discussion of the Kitaev table

Gian Michele Graf, ETH Zurich

PhD School: September 16-20, 2019
@Università degli Studi Roma Tre

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based on discussions with J. Haag, B. Roos

Outline

The periodic table of topological matter

Symmetry				d							
Class	Θ	Σ	Π	1	2	3	4	5	6	7	8
A	0	0	0	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}
AIII	0	0	1	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0
AI	1	0	0	0	0	0	\mathbb{Z}	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
BDI	1	1	1	\mathbb{Z}	0	0	0	\mathbb{Z}	0	\mathbb{Z}_2	\mathbb{Z}_2
D	0	1	0	\mathbb{Z}_2	\mathbb{Z}	0	0	0	\mathbb{Z}	0	\mathbb{Z}_2
DIII	-1	1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0	\mathbb{Z}	0
AII	-1	0	0	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0	\mathbb{Z}
CII	-1	-1	1	\mathbb{Z}	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0
C	0	-1	0	0	\mathbb{Z}	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0
CI	1	-1	1	0	0	\mathbb{Z}	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0

Notation for symmetries:

- ▶ Θ (time-reversal): antiunitary, $H\Theta = \Theta H$, $\Theta^2 = \pm 1$
- ▶ Σ (charge-conjugation): antiunitary, $H\Sigma = -\Sigma H$, $\Sigma^2 = \pm 1$
- ▶ $\Pi = \Theta\Sigma = \Sigma\Theta$: unitary

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First version: Schnyder et al.; then Kitaev based on
Altland-Zirnbauer; based on Bloch theory

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By now: Non-commutative (bulk) index formulae have been found in all cases (Prodan, Schulz-Baldes)

Basic symmetries

V : complex Euclidean vector space

Symmetries: (anti-)unitary maps $U : V \rightarrow V$ (names conventional)

1. Time-reversal $\Theta : V \rightarrow V$ antiunitary, $\Theta^2 =: \alpha = \pm 1$

If $V = V_+ \oplus V_-$:

2. Particle-hole $\Sigma : V_{\pm} \rightarrow V_{\mp}$ antiunitary, $\Sigma^2 =: \beta = \pm 1$
3. Chiral symmetry $\Pi : V_{\pm} \rightarrow V_{\mp}$ unitary

Remarks. 1) Let $U : V \rightarrow V$ with $U^2 = \gamma$; consider $\tilde{U} := cU$ with $c \in \mathbb{C}$, $|c| = 1$ to be chosen. Then

- ▶ U unitary: $\tilde{U}^2 = c^2\gamma$. W.l.o.g. γ arbitrary.
- ▶ U antiunitary: By $U^2U = UU^2$ we have $\gamma = \bar{\gamma}$, i.e. $\gamma = \pm 1$.
Intrinsic γ : $\tilde{U}^2 = |c|^2\gamma = \gamma$

2) For items 2, 3: $\dim V_+ = \dim V_-$.

3) Σ qualifies as Θ , so far.

Combination of symmetries

$$V = V_+ \oplus V_-$$

$$1'. \Theta : V_{\pm} \rightarrow V_{\pm}$$

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$$2. \Sigma : V_{\pm} \rightarrow V_{\mp}$$

$$3. \Pi : V_{\pm} \rightarrow V_{\mp}$$

Let

$$\Pi = \Theta\Sigma, \quad [\Theta, \Sigma] = 0$$

Then any two symmetries imply the third; moreover $\Pi^2 = \alpha\beta = \pm 1$

Remarks. 1) $\Sigma \neq \Theta$ (flip/no flip)

2) $\Pi \neq \Sigma, \Theta$ (unitary/antiunitary)

3) Possible combinations (none, one, three): $1 + 5 + 4 = 10$

symmetry classes

The classification

- ▶ Each entry of the table shows a group $G = 0, \mathbb{Z}, \mathbb{Z}_2$ (index group)
- ▶ Vector bundles over \mathbb{T}^d (torus) of a given symmetry class (“topological insulators”) are assigned an index $I \in G$
- ▶ If two of them (with indices I, I') are homotopy equivalent (within the class), then $I = I'$ (strong index).
- ▶ However, this is true only if their restrictions to all tori $\mathbb{T}^{d'} \subset \mathbb{T}^d$, ($d' < d$) are homotopy equivalent (weak indices).
- ▶ However, also non homotopy equivalent bundles may have $I = I'$, if they are so upon addition of trivial ones (stably homotopic \rightarrow K-theory)

The K-theoretic point of view

Example: Integers $k \in \mathbb{Z}$ may be identified with pairs $(n_+, n_-) \in \mathbb{N}^2$ of naturals, up to equivalence $(n'_+, n'_-) \sim (n''_+, n''_-)$ defined by

$$(n'_+ + \tilde{n}, n'_- + \tilde{n}) = (n''_+ + \tilde{\tilde{n}}, n''_- + \tilde{\tilde{n}})$$

for some $\tilde{n}, \tilde{\tilde{n}} \in \mathbb{N}$.

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Say $V' \sim V''$ if

$$(V'_+ \oplus \tilde{V}, V'_- \oplus \tilde{V}) \cong (V''_+ \oplus \tilde{\tilde{V}}, V''_- \oplus \tilde{\tilde{V}})$$

(homotopy).

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Say $V' \sim V''$ if

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(homotopy). In case of symmetry $U = \Theta, \Sigma, \Pi$, the map $U : V_{\pm} \rightarrow V_{\pm/\mp}$ is augmented to $U \oplus \tilde{U}$ with $\tilde{U} : \tilde{V} \rightarrow \tilde{V}$, $\tilde{U}^2 = \pm 1$

Remarks. 1) Dimension redefined (only here):

$$\dim V = \dim V_+ - \dim V_-$$

2) If Σ, Π are symmetries, $\dim V = 0$.

3) Notions extended to vector bundles

The derivation of the table

One more column $d = 0$: Vector bundles over a point \equiv vector spaces

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(Same as claimed for $d = 8$).

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(Same as claimed for $d = 8$). We'll derive it.

Normal forms of lone symmetries

Let

$$\varepsilon = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad \omega = i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\det \varepsilon = \det \omega = 1$$

We'll see: Depending on **classes**, $N := \dim V = n, 2n, 4n$,
($n = 1, 2, \dots$).

We'll construct **adapted** orthonormal bases $\underline{V} = (v_j)_{j=1}^k = (v_1, \dots, v_k)$
of invariant subspaces of V of dimension $k = 1, 2, 4$; generated by
arbitrary v_1 , ($\|v_1\| = 1$).

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Left action: $\underline{UV} = (Uv_1, \dots, Uv_k)$ ($U: V \rightarrow V$ map)

Right action: $\underline{VM} = (\sum_i v_i M_{ij})_{j=1}^k$ (M : matrix of order k)

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- ▶ Θ ($\alpha = +1$) has $k = 1$: v_1 such that $\Theta v_1 = v_1$, i.e.

$$\Theta \underline{V} = \underline{V}$$

- ▶ Θ ($\alpha = -1$) has $k = 2$: $v_2 := \Theta v_1$, $\Theta v_2 = -v_1$; then

$$\Theta \underline{V} = (\Theta v_1, \Theta v_2) = (v_2, -v_1) = \underline{V} \varepsilon$$

Normal forms of lone symmetries (cont.)

- ▶ $\Sigma (\beta = +1)$ has $k = 2$:

Normal forms of lone symmetries (cont.)

- ▶ Σ ($\beta = +1$) has $k = 2$: $v_1 \in V_+$, $v_2 := -i\Sigma v_1 \in V_-$; then

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- ▶ $\Pi (\gamma = -1 \text{ w.l.o.g.})$ has $k = 2$:

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- ▶ Σ ($\beta = +1$) has $k = 2$: $v_1 \in V_+$, $v_2 := -i\Sigma v_1 \in V_-$; then

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- ▶ Π ($\gamma = -1$ w.l.o.g.) has $k = 2$: $v_1 \in V_+$, $v_2 := -i\Pi v_1 \in V_-$; then

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Normal forms of combined symmetries Θ, Σ

- ▶ $\alpha = +1, \beta = \pm 1$ has $k = 2$:

$$\Theta \underline{V} = \underline{V}, \quad \Sigma \underline{V} = \underline{V} \begin{cases} \omega & (\beta = +1) \\ \varepsilon & (\beta = -1) \end{cases}$$

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- ▶ $\alpha = -1, \beta = \pm 1$ has $k = 4$:

$$\Theta \underline{V} = \underline{V} \begin{pmatrix} \varepsilon & 0 \\ 0 & \varepsilon \end{pmatrix}, \quad \Sigma \underline{V} = \underline{V} \begin{pmatrix} 0 & \beta 1_2 \\ 1_2 & 0 \end{pmatrix}$$

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Let $K : \mathbb{C}^N \rightarrow \mathbb{C}^N$ be the standard complex conjugation

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Claim: **Lone** symmetries $U = \Theta, \Sigma$ ($U^2 = \gamma = \pm 1$; $\gamma = \alpha$ **or** β , hence 2 + 2 cases) on V can be brought to the following form on \mathbb{C}^N

- ▶ $\gamma = +1$ ($N = n, 2n$): $U = K$
- ▶ $\gamma = -1$ ($N = 2n$): $U = \varepsilon K$

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2) For $U = \Sigma$ the split $V = V_+ \oplus V_-$ is compatibly realized as

$$V_{\pm} = \{(v, \pm iv) \mid v \in \mathbb{C}^n\} \subset \mathbb{C}^n \oplus \mathbb{C}^n = \mathbb{C}^N, \quad (\text{or flipped})$$

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4) This foreshadows: The lone symmetry Θ will not contribute to G (to be checked).

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- ▶ Classification: Right cosets \mathcal{B}/\mathcal{T} , connected components thereof.

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Remarks. 1) The index trivializes for Σ, Π by $\dim V_+ - \dim V_- = 0$.
Do new indices appear?

2) The index survives for just Θ (classes AI, AII). Does the group become larger? (Likely not by earlier remark)

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$$\varepsilon T = T \varepsilon, \quad \text{i.e. } T_- = T_+$$

- ▶ So $\mathcal{T} = U(n)$
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If further symmetries: $[H, \Theta] = 0$.

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$$A = UNU^T$$

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$U(n)$ is connected, hence index group $G = 0$.

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$$H = \begin{pmatrix} 0 & U^* \\ U & 0 \end{pmatrix}, \quad (U \in U(n)), \quad \Theta = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} K$$

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- ▶ $G = \pi_0(O(n)) = \mathbb{Z}_2$ by $\det U = \pm 1$

Class CI: $\alpha = +1, \beta = -1$

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- ▶ So $U\varepsilon U^T = \varepsilon$, i.e.
- ▶ $U \in Sp(2n) \cap U(2n)$, which is a connected group.
- ▶ Hence $G = 0$

Conclusion

Class	Symmetry			d
	Θ	Σ	Π	0
A	0	0	0	\mathbb{Z}
AIII	0	0	1	0
AI	1	0	0	\mathbb{Z}
BDI	1	1	1	\mathbb{Z}_2
D	0	1	0	\mathbb{Z}_2
DIII	-1	1	1	0
AII	-1	0	0	\mathbb{Z}
CII	-1	-1	1	0
C	0	-1	0	0
CI	1	-1	1	0