

Magnetic Weyl Quantization and Semiclassical Limit

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2019. 9. 19.

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What is Semiclassical Limit?

Semiclassical Limit

- Approximate quantum mechanics using classical mechanics.
- Introducing a semiclassical parameter $\varepsilon > 0$.
- When the limit of $\varepsilon \rightarrow 0$ is taken, we expect classical mechanics and quantum mechanics to match.

$$\left\{ \begin{array}{l} \frac{d}{dt}x = \nabla_{\xi}h \\ \frac{d}{dt}\xi = -\nabla_x h + B\dot{x} \end{array} \right. \quad \overset{?}{\longleftrightarrow} \quad \left\{ \begin{array}{l} i\varepsilon \frac{\partial \psi(t)}{\partial t} = \hat{H}^A \psi(t) \\ \psi(0) = \psi_0 \in L^2(\mathbb{R}^d) \end{array} \right.$$

f : Classical Observable, i.e. real function on $\mathbb{R}^d \times \mathbb{R}^d$

B : Magnetic Field, A : Vector Potential of Magnetic Field,

When $d = 2$, A satisfies $\text{rot}A = B$.

Motivation

Motivation

- To approximate time evolution generated by Quantum Hamiltonian \hat{H}^A with time evolution generated by Classical Hamiltonian h .

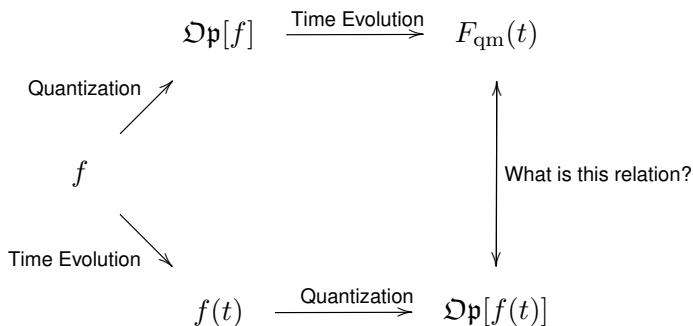


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1 Introduction

2 **Semiclassical Limit**

- Definition of Magnetic Weyl Quantization
- Time Evolution in Classical System and Quantum System
- Egorov type Theorem

Magnetic Weyl Quantization (Măntoiu, Purice 2004)

Magnetic Weyl Quantization

We consider integral of f :

$$\mathfrak{Op}^A[f] := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d \times \mathbb{R}^d} \mathcal{F}_\sigma[f](Y) W_\varepsilon^A(Y) dY.$$

We call a map $f \mapsto \mathfrak{Op}^A[f]$ as Magnetic Weyl Quantization of f .

$$\sigma((x, \xi), (y, \eta)) := \xi \cdot y - x \cdot \eta, \quad (\text{syplectic form}),$$

$$W_\varepsilon^A(Y) := \exp(i\sigma(Y, (Q, P_\varepsilon^A))), \quad (Y = (y, \eta) \in \mathbb{R}^d \times \mathbb{R}^d),$$

$$\mathcal{F}_\sigma[f](Y) := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d \times \mathbb{R}^d} e^{i\sigma(Y, Y')} f(Y') dY'.$$

Classical Time Evolution

$$\begin{aligned}
 f(t) &= f(x(t), \xi(t)), \\
 \frac{d}{dt}(f(t)) &= \sum_{j=1}^d \frac{\partial f}{\partial \xi_j} \frac{\partial h}{\partial x_j} - \frac{\partial f}{\partial x_j} \frac{\partial h}{\partial \xi_j} + \sum_{j,k=1}^d B_{jk} \frac{\partial f}{\partial \xi_j} \frac{\partial h}{\partial \xi_k} \\
 &=: \{h, f(t)\}_B, \quad f(0) = f
 \end{aligned}$$

Quantum time evolution

- $F_{\text{qm}}^A(t)$: Quantum observable at time $t \in \mathbb{R}$.

$$\begin{cases}
 i\varepsilon \frac{\partial}{\partial t} F_{\text{qm}}^A(t) = [F_{\text{qm}}^A(t), \hat{H}^A], & F_{\text{qm}}^A(0) = \mathfrak{Op}^A[f] \\
 F_{\text{qm}}^A(t) = e^{\frac{i}{\varepsilon} t \hat{H}^A} \mathfrak{Op}^A[f] e^{-\frac{i}{\varepsilon} t \hat{H}^A}.
 \end{cases}$$

Semiclassical Limit

Egorov Type Theorem(Lein [6] 2010)

$$F_{\text{cl}}(t) := \mathfrak{Op}^A[f(t)], \quad F_{\text{qm}}(t) := e^{i\frac{t}{\varepsilon}\hat{H}^A} \mathfrak{Op}^A[f] e^{-i\frac{t}{\varepsilon}\hat{H}^A},$$

Then there exists $C_T > 0$ such that,

$$\|F_{\text{cl}}(t) - F_{\text{qm}}(t)\|_{B(L^2(\mathbb{R}^d))} \leq \varepsilon^2 C_T, \quad (\forall t \in [-T, T])$$

Egorov type theorem(O. 2018)

There exists g_t such that

$$F_{\text{qm}}(t) = \mathfrak{Op}^A[g_t] + O_{B(L^2(\mathbb{R}^d))}(\varepsilon^\infty) \quad (1)$$

$$|\partial_\xi^\alpha \partial_x^\beta (g_t(x, \xi) - f(x(t), \xi(t)))| \leq C_T \varepsilon^2 \quad (2)$$

$$g_t(x, \xi) \asymp \sum_{n=0}^{\infty} \varepsilon^{2n} g_n(x, \xi).$$







Summary

- Magnetic Weyl Quantization is good quantization when magnetic field exists :





$$\mathfrak{Op}^A[f] := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d \times \mathbb{R}^d} \mathcal{F}_\sigma[f](Y) W_\varepsilon^A(Y) dY.$$

- The difference in time evolution generated by Classical Hamiltonian h and Quantum Hamiltonian $\mathfrak{Op}^A[h]$ is small.

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