

# Spectral statistics and many-body localization

Jan Šuntajs, Lev Vidmar, Janez Bonča

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Jožef Stefan Institute, Ljubljana, Slovenia

# Many-body localization (MBL) - what is it about?

- 1 Occurring in **INTERACTING** quantum systems with **DISORDER**
- 2 An **IDEAL INSULATOR** → at **ANY** temperature
- 3 Explains the **FAILURE** of some systems to **THERMALIZE**

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# The outline of the presentation

- 1 The properties of MBL systems
- 2 Introduction of the physical model
- 3 Numerical analysis of the **spectral statistics**
  - A brief introduction to spectral statistics
  - **Spectral form factor (SFF)**
- 4 Our recent results and conclusion

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- 4 Quantum chaos challenges many-body localization

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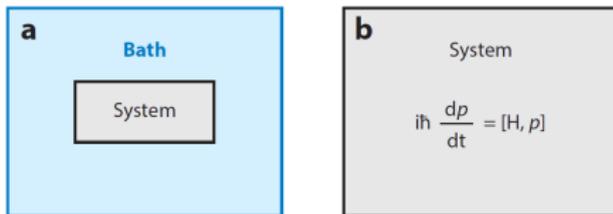
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- **Closed** quantum systems

Nandkishore, Huse, 2015

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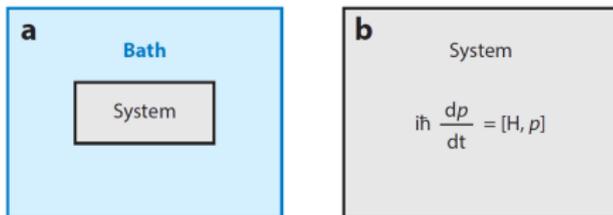


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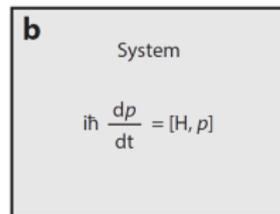
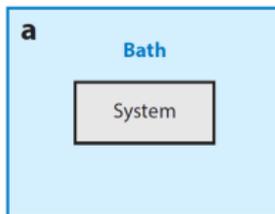


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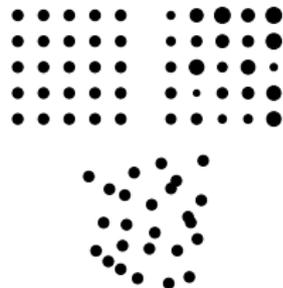
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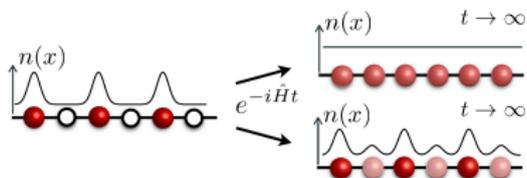
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- The **absence of ergodicity**

Abanin, Altman, Bloch, Serbyn, 2018

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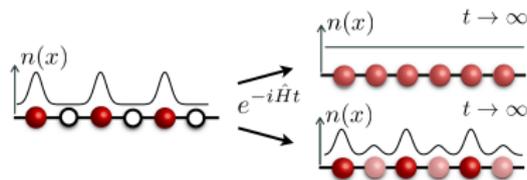


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- THE ENTANGLEMENT ENTROPY:
  - Area law scaling for all eigenstates
  - Logarithmic growth in time

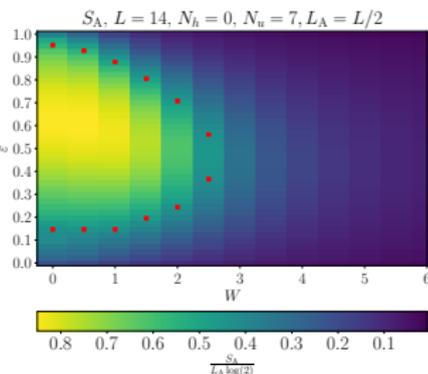
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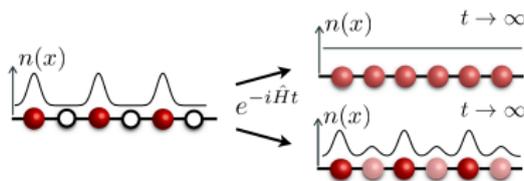
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PHYSICAL REVIEW B 77, 064426 (2008)

## Many-body localization in the Heisenberg $XXZ$ magnet in a random field

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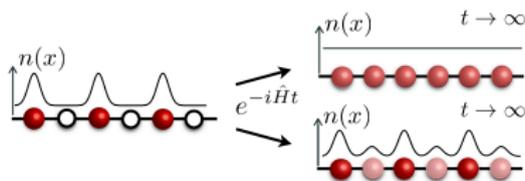
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(Received 31 August 2007; revised manuscript received 8 November 2007; published 25 February 2008)

- **Special properties of the energy spectra**
  - The subject of our numerical analysis

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# The model

The hamiltonian - **paradigmatic** 'quantum chaotic'/ergodic:

Heisenberg-like

1D

$$\hat{H} = \sum_{j \in \{1,2\}} J_j \sum_{\ell}^L \left( \hat{s}_{\ell}^x \hat{s}_{\ell+j}^x + \hat{s}_{\ell}^y \hat{s}_{\ell+j}^y + \Delta_j \hat{s}_{\ell}^z \hat{s}_{\ell+j}^z \right) + \sum_{\ell=1}^L w_{\ell} \hat{s}_{\ell}^z$$

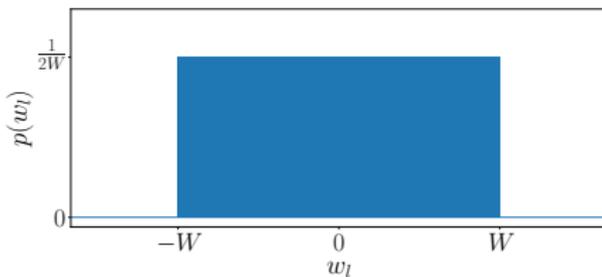
- $l$  - site,  $L$  - chain length
- $w_{\ell}$ : randomly disordered potential

# The model

$$J_1 = J_2 = 1$$

$$\Delta_1 = \Delta_2 = 0.55$$

Disorder probability distribution:



$W$  - the disorder strength parameter

# Our (numerical) analysis of the MBL systems

- We perform **full** or **partial** diagonalization of the Hamiltonians  
**Spectrum:**  $\{E_1 \leq E_2 \leq \dots \leq E_D\}$
- **Partial diagonalization:**  $\approx 500$  eigenstates from the middle of the spectra
  - Maximum Hilbert space dimensions:  
 $D = 48620$  (full)  
 $D = 184756$  (partial)
- Between  $10^2 - 10^3$  **different** disorder realizations for **each model parameter**

# Quantum chaos and energy spectra

## Why do we study energy spectra?

- Quantum chaos conjecture (Bohigas, Giannoni, Schmidt, 1984):

### Quantum systems

**Spectral properties** match the predictions of the **random matrix theory (RMT)**.

### Corresponding classical systems

The dynamics are completely **chaotic**.

- **What about systems without a classical analogue?**

# Quantum chaos and energy spectra

- **Many-body quantum chaos**

**GENERIC** systems → **RMT-like** spectral statistics

Montambaux *et. al.* (1993), Prosen (1999), Santos and Rigol (2008)

**RMT** statistics → hallmarks of **ergodicity** and **thermalization** in an isolated quantum system.

D'Alessio, Kafri, Polkovnikov, Rigol (2016)

# Statistical properties of the energy spectra

- We analyse the **statistical properties** of the energy spectra
- We rely on the findings of the **RMT**:
  - **Ergodic** systems: spectral statistic match the **Gaussian orthogonal ensemble (GOE)**
  - **MBL** systems: nearest levels distributed in accordance with the **Poisson** distribution
- We compare our **RESULTS** with the **above cases**.

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# The mean ratio of the level spacings

- The spacings between the **nearest** energy levels:

$$\delta_n = E_{n+1} - E_n \geq 0$$

- We define the **level spacing ratio**:

$$0 \leq \tilde{r}_n = \min\{\delta_n, \delta_{n-1}\} / \max\{\delta_n, \delta_{n-1}\} \leq 1$$

- **KEYNOTE:** the limiting values of  $\langle \tilde{r} \rangle$  are well known:

$$\text{Ergodic: } \langle \tilde{r} \rangle_{\text{GOE}} = 0.5307$$

$$\text{MBL: } \langle \tilde{r} \rangle_{\text{P}} = 2 \ln 2 - 1 \approx 0.3863$$

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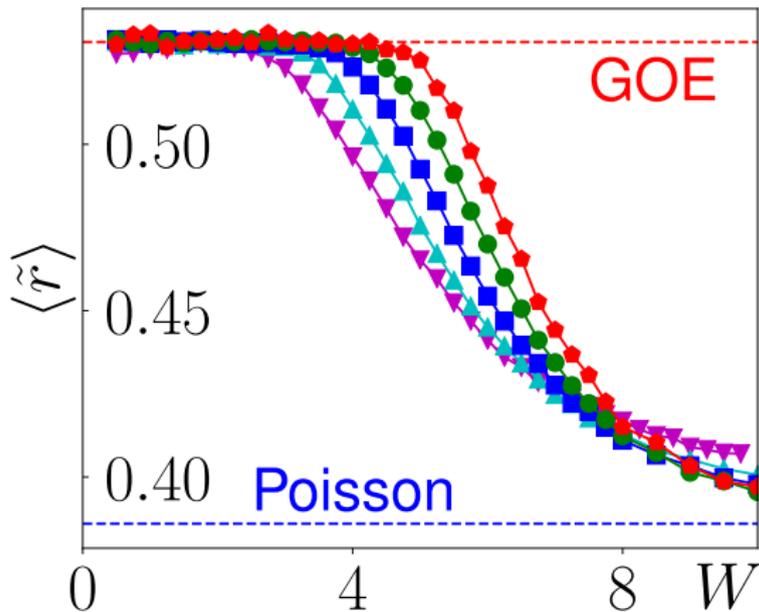
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# The mean ratio of the level spacings

▼  $L = 12$  ▲  $L = 14$  ■  $L = 16$  ●  $L = 18$  ◆  $L = 20$



## Mean level spacings ratio:

- a commonly used indicator of a given system's ergodicity
  - + straightforward implementation
- - only considers correlation between the **nearest** energy levels

We would like to consider correlations between all the levels

- This is why we implement the **spectral form factor (SFF)**
  - the implementation is more demanding

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# The spectral form factor (SFF)

Definition:

$$K(\tau) := \left\langle \frac{1}{D} \sum_{i,j}^D e^{-i(\varepsilon_i - \varepsilon_j)\tau} \right\rangle; \quad K(0) = D, \quad K(\tau \rightarrow \tau_H) = 1$$

$D$  - Hilbert space dimension       $\tau \rightarrow$  an external parameter

$\langle \dots \rangle$  over disorder realizations

**Heisenberg** time  $\tau_H \propto$  inverse *mean level spacing*  
(largest **sensible** timescale of a system)

$\{\varepsilon_i\} \rightarrow$  energy levels after *spectral unfolding*

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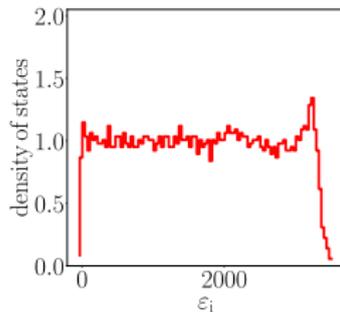
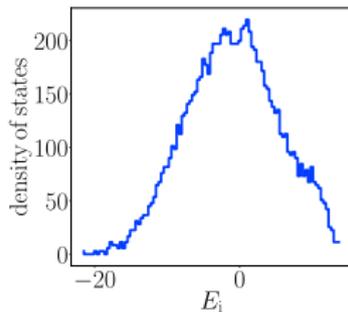
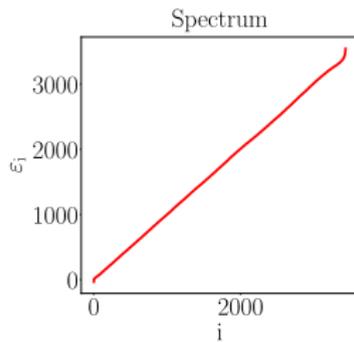
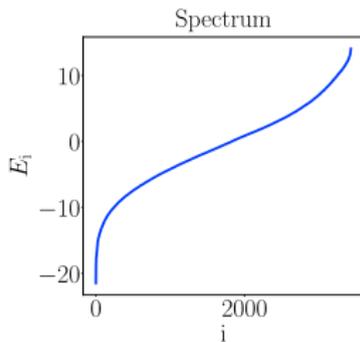
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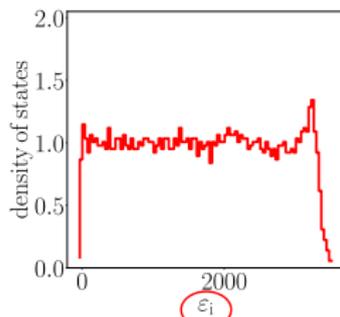
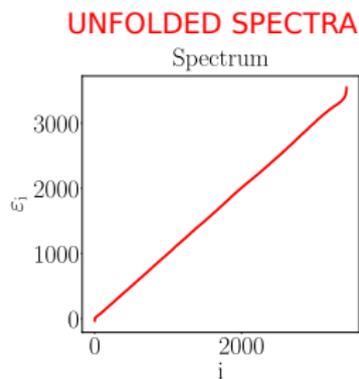
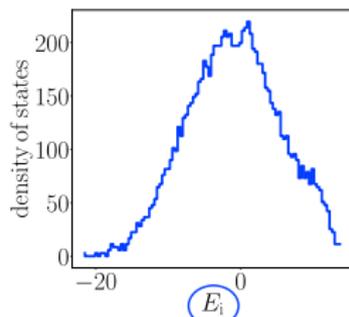
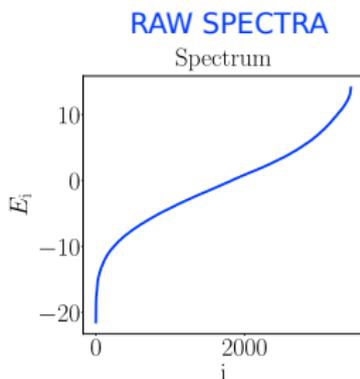
# The spectral form factor (SFF)

- A quick introduction to unfolding



# The spectral form factor (SFF)

- **Unfolding** → mean level spacing = 1

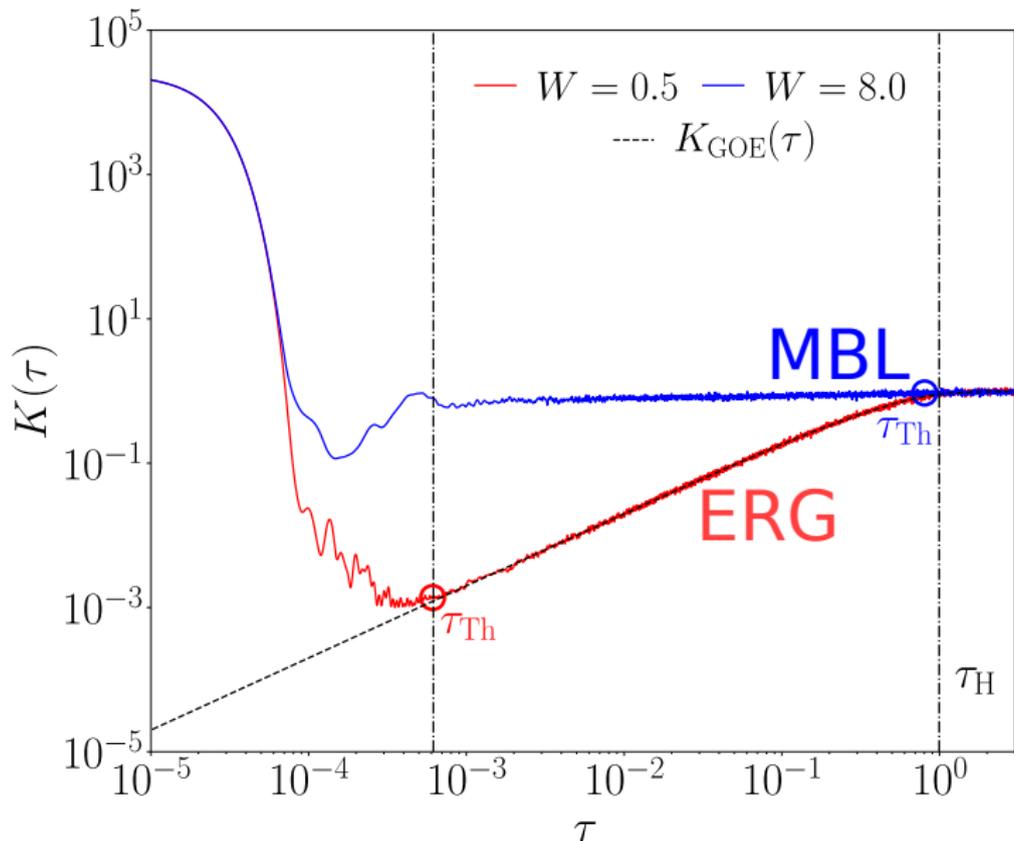


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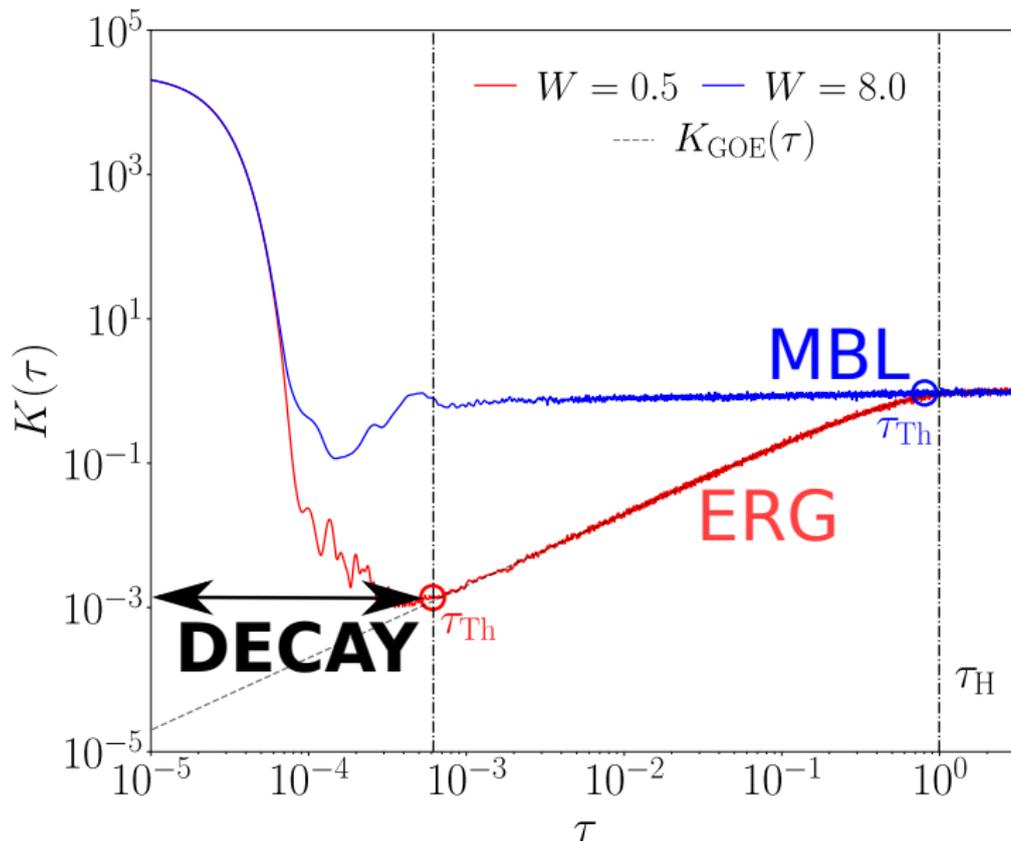
## SFF - KEYNOTES:

- we again expect different behaviour for **ergodic** and **uncorrelated** spectra
- we investigated the behaviour of the **Thouless time**  $\tau_{\text{Th}}$
- **Thouless time**  $\tau_{\text{Th}}$   $\rightarrow$  the onset of **UNIVERSAL DYNAMICS**

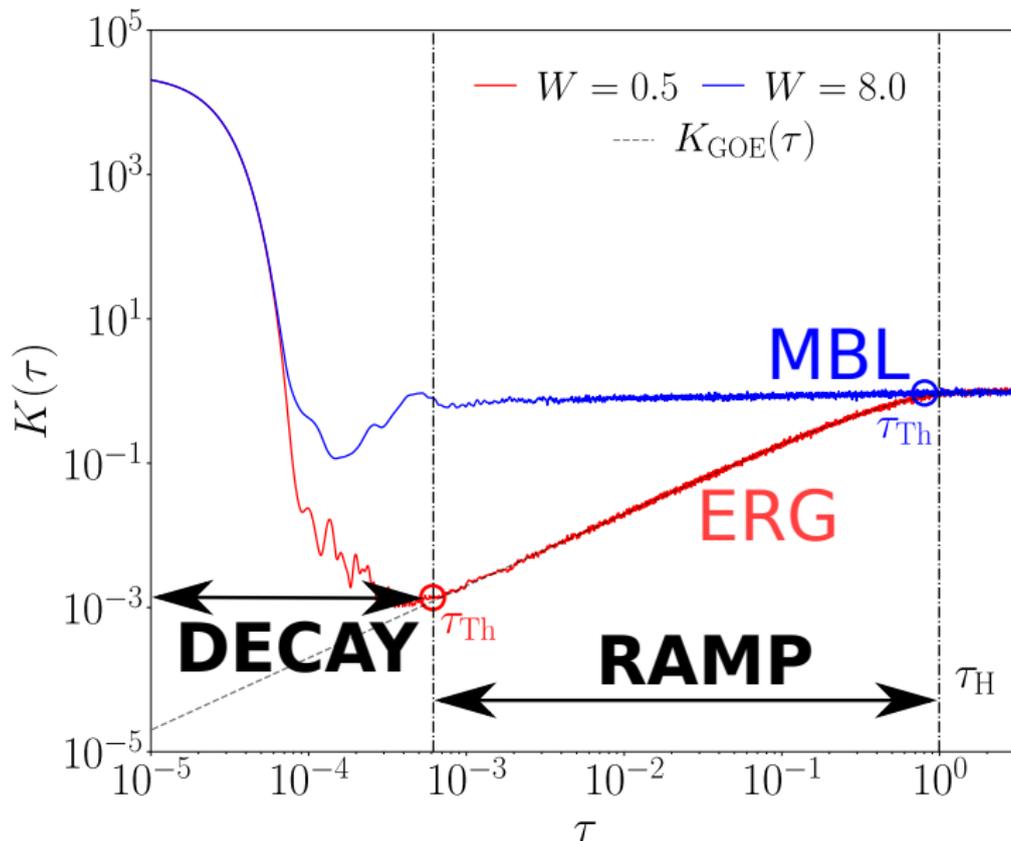
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We see universal behaviour after some transient time  $\tau_{\text{Th}}$

$K(\tau)$  for **uncorrelated** spectra

$$K_{\text{P}}(\tau) = 1$$

$K(\tau)$  in ergodic systems

$$K_{\text{GOE}}(\tau) = 2 - \tau \log \left( \frac{2\tau + 1}{2\tau - 1} \right)$$

# SFF - explaining the Thouless time $\tau_{\text{Th}}$

- **Thouless time**  $\tau_{\text{Th}}$ : determines the energy scale at which the **spectral correlations** are universally determined by the **GOE** predictions (e.g., when the **RAMP** appears)
  
- large(r)  $\tau_{\text{Th}}$   $\rightarrow$  small(er) **spectral correlation length**  $E_{\text{Th}}$

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- large(r)  $\tau_{\text{Th}}$   $\rightarrow$  small(er) **spectral correlation length**  $E_{\text{Th}}$

- **We set out to find out different scalings**
- How does  $\tau_{\text{Th}}$  scale with the system size  $L$ ?
- How does  $\tau_{\text{Th}}$  scale with disorder strength parameter  $W$

## We noticed some surprising results along the way

Quantum chaos **challenges** many-body localization

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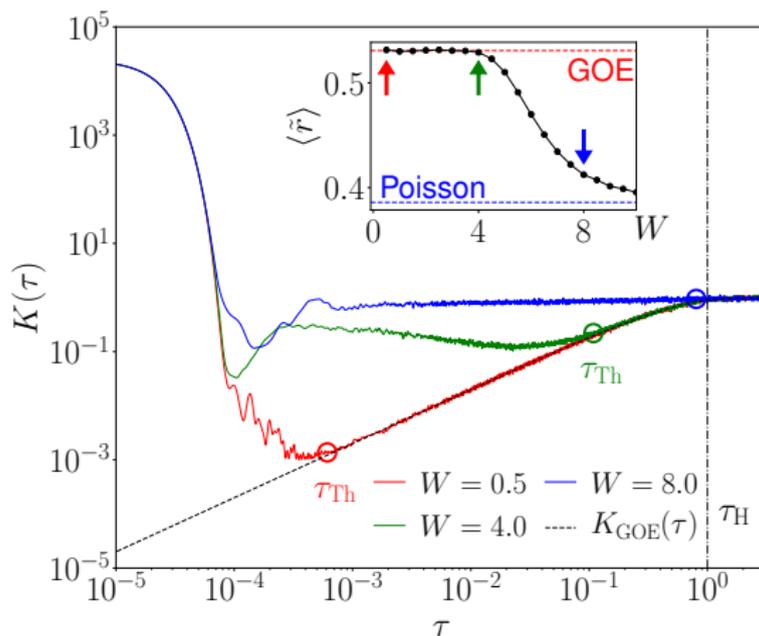
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# SFF - results

We first checked for consistency of the  $K(\tau)$  and  $\langle \tilde{r} \rangle$  results



$L = 18$

# SFF - results

- We then numerically extracted  $\tau_{\text{Th}}$  values and performed a scaling analysis w.r.t. both  $L$  and  $W$
- In the subsequent scaling analysis, we introduce the **PHYSICAL THOULESS TIME**  $t_{\text{Th}}$ , rescaling  $\tau_{\text{Th}}$  by the mean level spacing  $\overline{\delta E}$  of the **RAW** spectra:

$$t_{\text{Th}} = \tau_{\text{Th}} / \overline{\delta E}$$

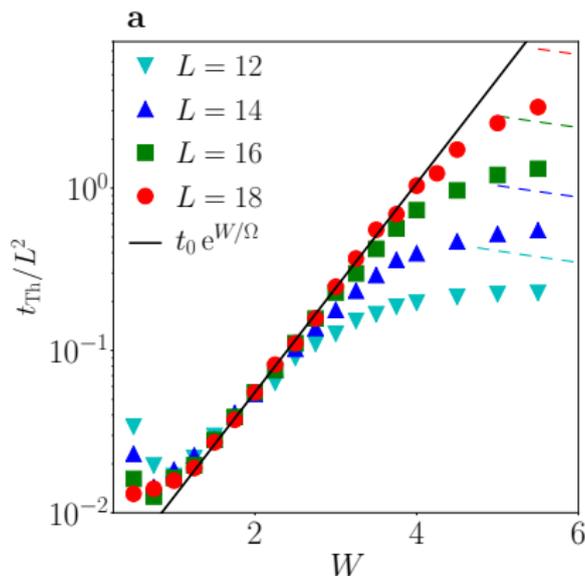
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The results of the scaling analysis

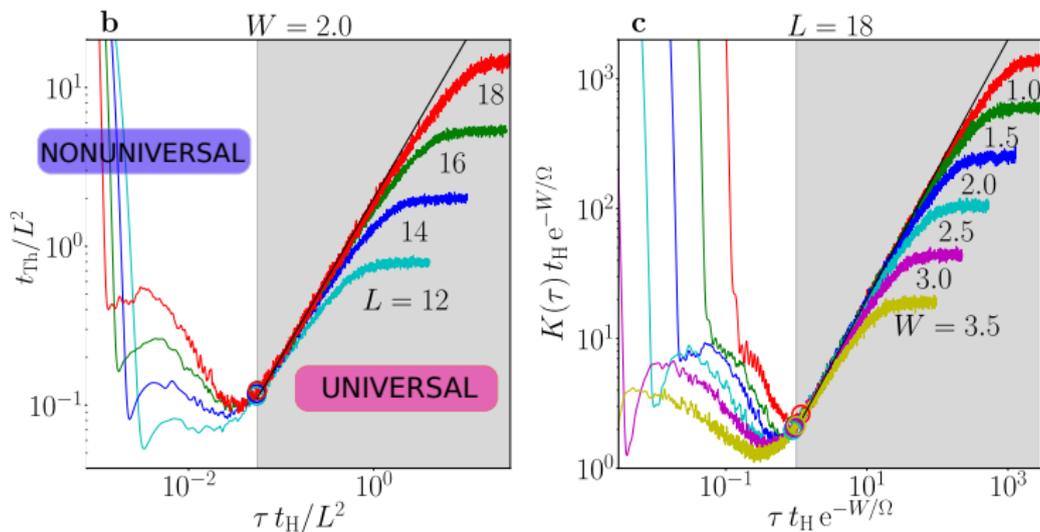


We obtain the following double scaling:

$$t_{\text{Th}} = t_0 e^{W/\Omega} L^2$$

# SFF - results

## Verification of the scaling: data collapse



# SFF - results

## ... Is there a MBL transition at all?

- a **SUPPOSED** MBL transition occurs for some **CRITICAL** disorder  $W^*$  when the energy spectrum becomes uncorrelated:

$$t_{\text{Th}}(W^*) = t_{\text{H}}$$

- $t_{\text{H}}$  scaling is given by:

$$t_{\text{H}} \propto \exp(L \ln 2)$$

- combining these results gives us

$$W^* \approx \Omega \ln(2) L \propto L$$

- **IMPLIES ABSENCE OF MBL IN THE THERMODYNAMIC LIMIT!**

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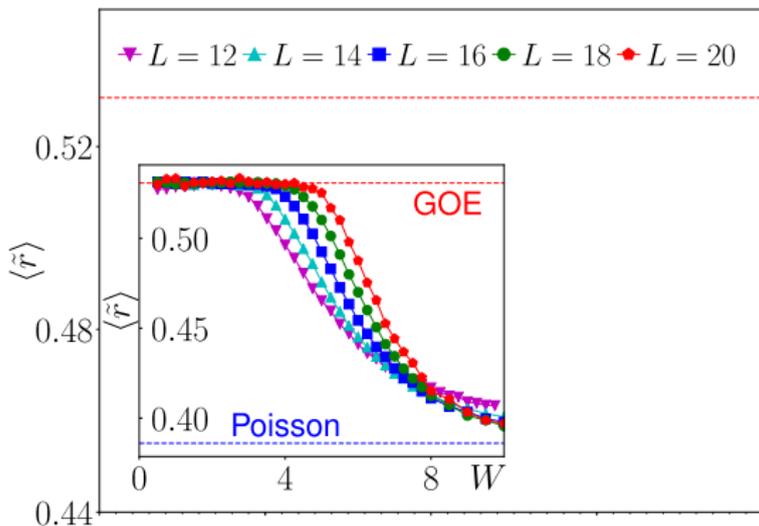
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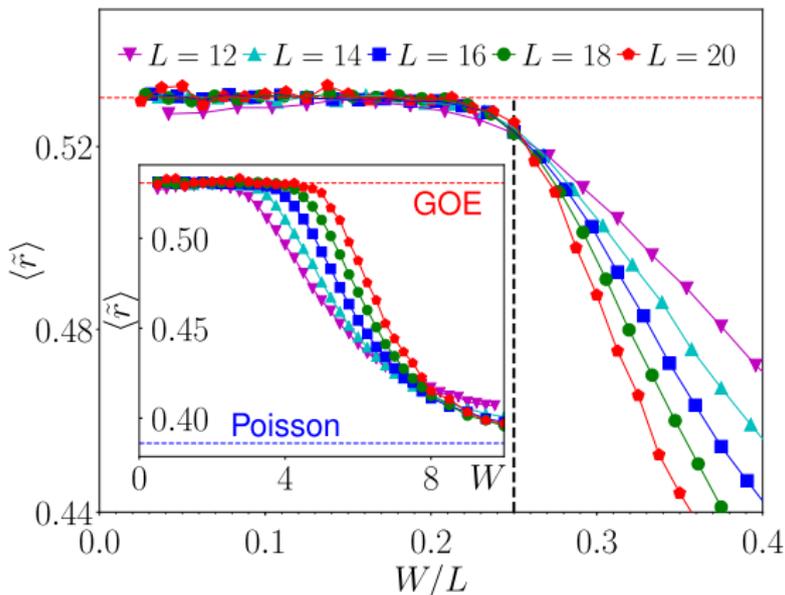
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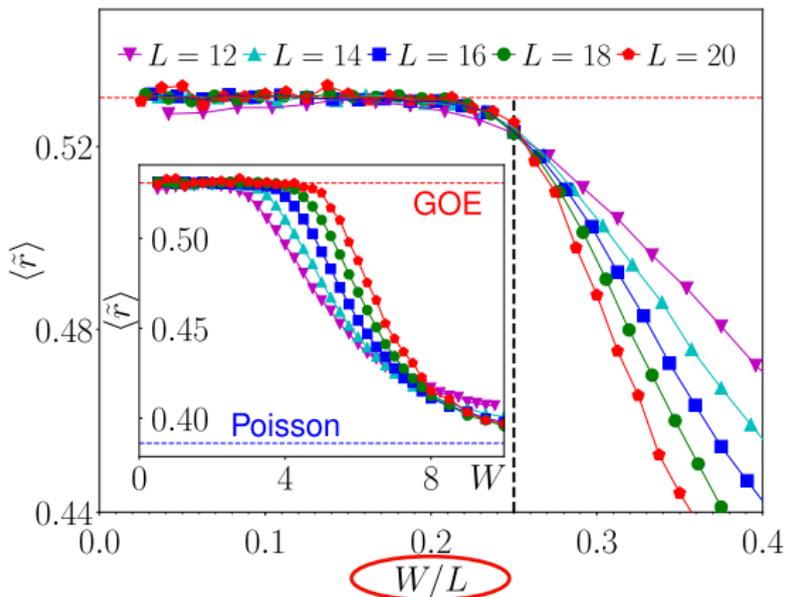
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# Verification of our results

- **BOTTOM LINE:** results for different spectral statistics seem consistent

# Conclusions and further work

- Our results for the **PARADIGMATIC** class of models expected to give **MBL** show **NO INDICATIONS** of the MBL transition
- The emergence of **QUANTUM CHAOS** for **ANY** disorder strength in the TD limit

# Conclusions and further work

- We need to test our assumptions on other models in which MBL is predicted
- We need to examine and better understand the relationship between **our results** and the **transport properties**

**Thank you for your attention!**