MS410 - Primo esonero (28-4-2025)

1. Consider the 1D Ising model with nearest neighbour interactions in a box of side L with periodic boundary conditions, $H_{h,L}^{per}(\sigma)$. Compute

$$f_{\beta,h}(x-y) := \lim_{L \to \infty} \langle \sigma_x \sigma_y \rangle_{\beta,h,L}^{per},$$

where $\langle \cdot \rangle_{\beta,h,L}^{per}$ is the grand-canonical average with respect to $H_{h,L}^{per}(\sigma)$. Show that $f_{\beta,h}(x)$ converges exponentially to $[m(\beta,h)]^2$ as $|x| \to \infty$ (here $m(\beta,h) = \beta^{-1}\partial_h\psi(\beta,h)$ is the average magnetization), namely that

$$f_{\beta,h}(x) - [m(\beta,h)]^2 \sim Ce^{-\kappa|x|},\tag{1}$$

for suitable $C, \kappa > 0$. In particular, compute the rate $\kappa = \kappa(\beta, h)$.

- 2. Consider the 1D Ising model with even, finite range, interaction J(x-y) of range r (i.e., J(x) = 0 for |x| > r) in a box of side L with periodic boundary conditions. Prove that the grand-canonical partition function $Z_{\beta,h,L}^{per}$ can be written as $\text{Tr}M^L$ for a suitable $2^r \times 2^r$ transfer matrix M, which in general is not symmetric. Note that the elements of M are non-negative and show that there exists a positive integer p such that M^p has positive elements: find the smallest value of p for which such property holds.
- 3. [Perron-Frobenius theorem] Consider a $n \times n$ real, not necessarily symmetric, matrix A with positive elements.
 - Using Brouwer fixed-point theorem (Every continuous function from a convex compact subset K of a Euclidean space to K itself has a fixed point) show that A admits a right eigenvector with positive components. [Hint: apply Brouwer fixed-point theorem to the map **v** → [∑_{i=1}ⁿ(A**v**)_i]⁻¹ A**v** from the space of vectors **v** with non-negative components such that ∑_{i=1}ⁿ(**v**)_i = 1 to itself.] Denote by **v**^{*} such eigenvector, normalized so that ∑_i(**v**^{*})_i = 1, and by λ^{*} > 0 its eigenvalue. Moreover, show that A admits a left eigenvector **u**^{*} with positive components and the same eigenvalue λ^{*}, normalized so that **u**^{*} · **v**^{*} = 1. [Hint: apply again Brouwer fixed-point theorem to the map **u** → (λ^{*})⁻¹A^T**u** from the space of vectors **u** with non-negative components such that **u** · **v**^{*} = 1 to itself.]
 - Show that the eigenvalue λ^* is simple and that any other eigenvalue λ of A is smaller than λ^* in absolute value at least by a factor $1 e^{-2c}$, i.e.,

$$|\lambda| \le (1 - e^{-2c})\lambda^*,$$

with $e^{-c} := \min_{i,j,k} \frac{A_{ij}}{A_{kj}} \in (0, 1]$ the *Perron-Frobenius gap*. [Hint: let $A^* := A/\lambda^*$ and show that for any non-zero vector \mathbf{v} with nonnegative components $e^{-c} \leq \frac{(A^*\mathbf{v})_i}{(A^*\mathbf{v})_j} \leq e^c$ for any $1 \leq i, j \leq n$ and any $k \geq 1$; in particular $(\mathbf{v}^*)_i \geq e^{-c}(\mathbf{v}^*)_j$ for any i, j, from which $\max_j(\mathbf{v}^*)_j \sum_i (\mathbf{u}_i^*) \leq e^c$; using this one finds that, for any \mathbf{v} with non-negative components and any $i, (A^*\mathbf{v})_i \geq e^{-2c}(\mathbf{u}^* \cdot \mathbf{v})(\mathbf{v}^*)_i$; therefore, if, given \mathbf{v} a real vector, we denote by $|\mathbf{v}|$ the vector with components $|\mathbf{v}|_i := |(\mathbf{v})_i|$, letting $\mathbf{v}_{\pm} = (|\mathbf{v}| \pm \mathbf{v})/2$, we find that, for any real vector \mathbf{g} such that $\mathbf{u}^* \cdot \mathbf{g} = 0$,

$$A^* \mathbf{g} = A^* (\mathbf{g}_+ - e^{-2c} (\mathbf{u}^* \cdot \mathbf{g}_+) \mathbf{v}^*) - A^* (\mathbf{g}_- - e^{-2c} (\mathbf{u}^* \cdot \mathbf{g}_-) \mathbf{v}^*),$$

so that, for any $1 \leq i \leq n$,

$$|A^*\mathbf{g}|_i \le \left(A^*(\mathbf{g}_+ - e^{-2c}(\mathbf{u}^* \cdot \mathbf{g}_+)\mathbf{v}^*) + A^*(\mathbf{g}_- - e^{-2c}(\mathbf{u}^* \cdot \mathbf{g}_-)\mathbf{v}^*)\right)_i.$$

Hence, $\mathbf{u}^* \cdot |A^*\mathbf{g}| \leq (1 - e^{-2c})(\mathbf{u}^* \cdot |\mathbf{g}|)$ so that, iterating, $\mathbf{u}^* \cdot |(A^*)^k \mathbf{g}| \leq (1 - e^{-2c})^k (\mathbf{u}^* \cdot |\mathbf{g}|)$. Therefore, given any real vector \mathbf{v} , writing $\mathbf{v} = (\mathbf{u}^* \cdot \mathbf{v})\mathbf{v}^* + \mathbf{g}$ with $\mathbf{g} = \mathbf{v} - (\mathbf{u}^* \cdot \mathbf{v})\mathbf{v}^*$, we find $\mathbf{u}^* \cdot |(A^*)^k \mathbf{v} - (\mathbf{u}^* \cdot \mathbf{v})\mathbf{v}^*| \leq 2(1 - e^{-2c})^k \mathbf{u}^* \cdot |\mathbf{v}|$, from which the result follows.]

Consider now a matrix A with non-negative elements, such that there exists a positive integer $p \geq 1$ such that A^p has positive components. Note that by applying the previous results to A^p , one finds that also A admits a unique maximal eigenvalue (such that all other eigenvalues of A are strictly smaller in absolute value) with left and right eigenvectors whose components are all positive.

4. In the context of Problem 2, apply Perron-Frobenius theorem (in the form discussed at the end of the previous problem) to M to conclude that ψ(β, h) = log λ(β, h) with λ(β, h) the unique maximal eigenvalue of M. Using the theorem of analytic dependence of the simple eigenvalues of analytic matrices (Let A(ε) be a n × n matrix whose elements are analytic in ε in a neighborhood U ⊂ C of ε₀ ∈ C; if λ₀ is a simple eigenvalue of A(ε₀), then there exists a neighborhood V ⊂ U of ε₀ and an analytic function λ : V → C such that λ(ε) is a simple eigenvalue of A(ε) for all ε ∈ V and λ(ε₀) = λ₀) to conclude that ψ(β, h) is real-analytic in β, h for all β > 0 and h ∈ ℝ. Next, define f_{β,h}(x − y) as in Problem 1, and prove that Eq.(1) holds.

Next, define $f_{\beta,h}(x-y)$ as in Problem 1, and prove that Eq.(1) holds. Give a lower bound on κ in terms of the Perron-Frobenius gap of the transfer matrix. 5. [The solution of the Curie-Weiss model via the Hubbard-Stratonovich transformation] Using the Gaussian identity

$$e^{\frac{\alpha}{2}x^2} = \frac{1}{\sqrt{2\pi\alpha}} \int_{-\infty}^{+\infty} e^{-\frac{m^2}{2\alpha} + mx} \, dm,$$

valid for all $\alpha > 0$, prove that the grandcanonical partition function of the Curie-Weiss model with coupling $J_0 > 0$ can be rewritten as:

$$Z_{\beta,h,N}^{CW} = \sqrt{\frac{N\beta J_0}{2\pi}} \sum_{\sigma_1,...,\sigma_N=\pm} \int_{-\infty}^{+\infty} e^{-N\beta J_0 m^2/2 + \beta (J_0 m + h) \sum_{i=1}^N \sigma_i} dm$$
$$= \sqrt{\frac{N\beta J_0}{2\pi}} \int_{-\infty}^{+\infty} e^{-N\beta J m^2/2} (2\cosh(\beta (J_0 m + h)))^N dm$$

from which it follows that

$$\lim_{N \to \infty} \frac{1}{N} \log Z_{\beta,h,N}^{CW} = \max_{m \in \mathbb{R}} \left[-\beta J_0 m^2 / 2 + \log \cosh[\beta (J_0 m + h)] + \log 2 \right].$$

Verify explicitly that this expression is the same as the expression $\psi^{CW}(\beta, h) = \max_{-1 \le m \le 1} \{\beta hm - \beta f^{CW}(\beta, m)\}$ computed in class.

- 6. Let $\psi(\beta, h)$ be the pressure of the *d*-dimensional Ising model with ferromagnetic interaction J(x-y), and $\psi^{CW}(\beta, h)$ the pressure of the Curie-Weiss model with coupling $J_0 := \sum_{x \neq 0} J(x)$. Prove that, if $h \ge 0$, then $\psi(\beta, h) \ge \psi^{CW}(\beta, h)$, via the following steps.
 - (a) Prove that, for any $m \in \mathbb{R}$, $H_{h,\Lambda}^{per}(\sigma)$ can be re-written as $H_{h,\Lambda}^{per}(\sigma) = H_{h,\Lambda}^{per,0}(\sigma) + H_{h,\Lambda}^{per,1}(\sigma)$, where

$$H_{h,\Lambda}^{per,0}(\sigma) = \frac{J_0}{2}m^2|\Lambda| - (J_0m + h)\sum_{x \in \Lambda}\sigma_x,$$

$$H_{h,\Lambda}^{per,1}(\sigma) = -\frac{1}{2}\sum_{\substack{x,y \in \Lambda:\\x \neq y}}J(x-y)(\sigma_x - m)(\sigma_y - m)$$

(b) Using the previous rewriting, recognize that the grand-canonical partition function of $H_{h,\Lambda}^{per}(\sigma)$ can be rewritten as

$$Z_{\beta,h,\Lambda}^{per} = Z_{\beta,h,\Lambda}^{per,0} \langle e^{-\beta H_{h,\Lambda}^{per,1}} \rangle_{\beta,h,\Lambda}^{per,0}$$

where $Z_{\beta,h,\Lambda}^{per,0}$ is the grand-canonical partition function of $H_{h,\Lambda}^{per,0}(\sigma)$ and $\langle (\cdot) \rangle_{\beta,h,\Lambda}^{per,0}$ is the average with respect to the grand-canonical distribution associated with $H_{h,\Lambda}^{per,0}(\sigma)$ at inverse temperature β . (c) Use Jensen's inequality (stating that $\int \mu(dx)f(x) \ge f(\int \mu(dx)x)$ for any probability measure μ and any **convex** function f), to conclude that

$$Z_{\beta,h,\Lambda}^{per} \ge Z_{\beta,h,\Lambda}^{per,0} e^{-\beta \langle H_{h,\Lambda}^{per,1} \rangle_{\beta,h,\Lambda}^{per,0}}.$$

Compute the right side explicitly as a function of m. Show that, by fixing m to be the largest solution of $m = \tanh[\beta(J_0m + h)]$, one obtains, after having taken the thermodynamic limit,

$$\psi(\beta, h) \ge \psi^{CW}(\beta, h), \quad \forall h \ge 0,$$

as desired.

- 7. Consider the Curie-Weiss model $H_{h,N}^{CW}(\sigma)$ with coupling $J_0 > 0$ and magnetic field $h \ge 0$. Let $m^*(\beta, h)$ be the largest solution of $m = \tanh \beta (J_0 m + h)$.
 - (a) Prove that, if h > 0, or if h = 0 and $\beta < \beta_c$, then, for all $\epsilon > 0$,

$$\lim_{N \to \infty} \mathbb{P}_{\beta,h,N}(|m - m^*(\beta,h)| > N^{-1/2+\epsilon}) = 0,$$

where $\mathbb{P}_{\beta,h,N}$ is the probability with respect to the grand-canonical distribution associated with $H_{h,N}^{CW}(\sigma)$ at inverse temperature β .

(b) Similarly, prove that, if h = 0 and $\beta > \beta_c$, then, for all $\epsilon > 0$,

$$\lim_{N \to \infty} \mathbb{P}_{\beta,0,N}(||m| - m^*(\beta,0)| > N^{-1/2+\epsilon}) = 0.$$

(c) Finally, prove that, if h = 0 and $\beta = \beta_c$, then, for all $\epsilon > 0$,

$$\lim_{N \to \infty} \mathbb{P}_{\beta_c, 0, N}(|m| > N^{-1/4 + \epsilon}) = 0.$$